

# Bayes Theorem (Recap)

Adrian Bevan

email: [a.j.bevan@qmul.ac.uk](mailto:a.j.bevan@qmul.ac.uk)



# Overview

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- Bayes theorem is given by

$$P(B|A) = \frac{P(A|B)}{P(A)} P(B)$$

The probability you want to compute: The probability of the hypothesis B given the data A. This is sometimes called the posterior probability



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The probability of the data A, given the hypothesis B. This you can compute given a theory or model

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The probability of the data A given all possible hypotheses.

$$P(A) = \sum_j P(A|B_j)P(B_j)$$



## Example

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- ▶ What is the probability of being dealt an ace from a deck of 52 cards?
  - ▶ Two hypothetical outcomes: you are dealt an ace ( $B_0$ ), and you are not dealt an ace ( $B_1$ ).
  - ▶ Assume that the deck is unbiased and properly shuffled, so that there are 4 aces out of 52 randomly distributed in the pack of cards.

$$P(B_0) = 4/52$$

$$P(B_1) = 48/52$$

Is this choice of prior reasonable?

- ▶ Compute the probability of being dealt an ace, given that you have an unbiased

$$P(A|B_0) = 4/52$$

$$P(A|B_1) = 48/52$$

- ▶ Now we can compute the posterior probability



$$P(A) = \sum_j P(A|B_j)P(B_j)$$

$$\begin{aligned} P(A) &= (4/52) \times (4/52) + (48/52) \times (48/52) \\ &= (16 + 2304)/2704 \\ &= 0.8580(4 \text{ s.f.}) \end{aligned}$$

$$P(B|A) = \frac{P(A|B)}{P(A)} P(B)$$

$$\begin{aligned} P(A|B_0) &= \frac{(4/52) \times (4/52)}{0.8580} \\ &= 0.007(4 \text{ s.f.}) \end{aligned}$$

This result does not agree with a frequentist interpretation of the data.







## An alternative calculation: assume uniform priors

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$$\begin{aligned}P(A) &= \sum_j P(A|B_j)P(B_j) \\P(A) &= (4/52) \times (1/2) + (48/52) \times (1/2) \\&= (4 + 48)/104 \\&= 0.5\end{aligned}$$

$$P(B|A) = \frac{P(A|B)P(B)}{P(A)}$$

$$\begin{aligned}P(A|B_0) &= \frac{(4/52) \times (1/2)}{0.5} \\&= 0.077(3 \text{ d.p.})\end{aligned}$$

This result agrees with a frequentist interpretation of the data.



# Weighted Average (Recap)

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## Formula for two uncorrelated variables

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$$\begin{aligned}\bar{x} \pm \sigma_x &= \frac{x_1/\sigma_1^2 + x_2/\sigma_2^2}{1/\sigma_1^2 + 1/\sigma_2^2} \pm (1/\sigma_1^2 + 1/\sigma_2^2)^{-1/2}, \\ &= \frac{\sigma_2^2 x_1 + \sigma_1^2 x_2}{\sigma_1^2 + \sigma_2^2} \pm \left( \frac{\sigma_1^2 \sigma_2^2}{\sigma_1^2 + \sigma_2^2} \right)^{1/2}.\end{aligned}$$

- ▶ For example, let's just consider the total precision on the measurement of some observable  $X$ . If there are two ways to measure  $X$ , and these yield  $\sigma_1 = 5$  and  $\sigma_2 = 3$ , then the total error on the average is

$$\begin{aligned}\left( \frac{\sigma_1^2 \times \sigma_2^2}{\sigma_1^2 + \sigma_2^2} \right)^{1/2} &= \left( \frac{25 \times 9}{25 + 9} \right)^{1/2} \\ &= \sqrt{6.62} \simeq 2.57\end{aligned}$$



- 
- ▶ If one has a set of uncorrelated observables, then it is straightforward to compute the average in the same way for each observable in the set.
  - ▶ i.e. instead of an  $n$  dimensional problem, you have  $n$  lots of a one dimensional problem as in the previous example.
  - ▶ Unfortunately if you have a set of correlated observables from different measurements, then this is no longer the case and the problem becomes a little more complicated.



- ▶ A more complicated problem arises when the observables are correlated... In this case the covariance matrix between a set of  $M$  measured observables will play a role.

inverse of the covariance matrix  
for the  $j^{\text{th}}$  measurement

$$\bar{x} = \left[ \sum_{j=1}^M V_j^{-1} \right]^{-1} \cdot \left[ \sum_{j=1}^M V_j^{-1} x_j \right],$$
$$V = \left[ \sum_{j=1}^M V_j^{-1} \right]^{-1}.$$

The  $j^{\text{th}}$  measurement vector (the  $n$  correlated observables)

- ▶ The first factor is common both to the covariance matrix for the average, and the observable values for the average.



# Example: Measuring time-dependent CP asymmetries in B meson decays

- Real example using results from two HEP experiments

Experiment	$S$	$C$	$\rho_{SC}$	$\text{cov}_{SC}$
<i>BABAR</i> (Aubert <i>et al.</i> , 2007)	$-0.170 \pm 0.207$	$+0.010 \pm 0.162$	0.035	-0.0012
Belle (Somov <i>et al.</i> , 2007)	$+0.190 \pm 0.310$	$-0.160 \pm 0.225$	0.100	0.0070

$M = 2$

$n = 2$

Remember if you only have a correlation, you need to compute the covariance in order to compute  $V_j$ .



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$$V_1 = \begin{pmatrix} 0.0430 & -0.0012 \\ -0.0012 & 0.0261 \end{pmatrix},$$

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- ▶ Having constructed  $x_j$  and  $V_j$ , one can compute the average ...



- ▶ Starting with the error matrix:

$$V = \left[ \sum_{j=1}^M V_j^{-1} \right]^{-1}.$$

- ▶ Where

$$V_1^{-1} + V_2^{-1} = \begin{pmatrix} 33.749 & -0.402 \\ -0.402 & 58.363 \end{pmatrix},$$

- ▶ So

$$V = \begin{pmatrix} 0.0296 & 0.0002 \\ 0.0002 & 0.0171 \end{pmatrix},$$

**Results:**

Covariance = 0.0002

Variances are read off of the diagonal



- ▶ Similarly for the average value we can now compute S and C (the n observables of interest in our example) via:

$$\bar{x} = \left[ \sum_{j=1}^M V_j^{-1} \right]^{-1} \cdot \left[ \sum_{j=1}^M V_j^{-1} x_j \right],$$

- ▶ Thus the average of the two measurements is:

$$\begin{pmatrix} S \\ C \end{pmatrix} = \begin{pmatrix} -0.05 \pm 0.17 \\ -0.06 \pm 0.13 \end{pmatrix}$$

- ▶ with a covariance of 0.0002 between S and C.

# Poisson limits

Adrian Bevan

email: [a.j.bevan@qmul.ac.uk](mailto:a.j.bevan@qmul.ac.uk)



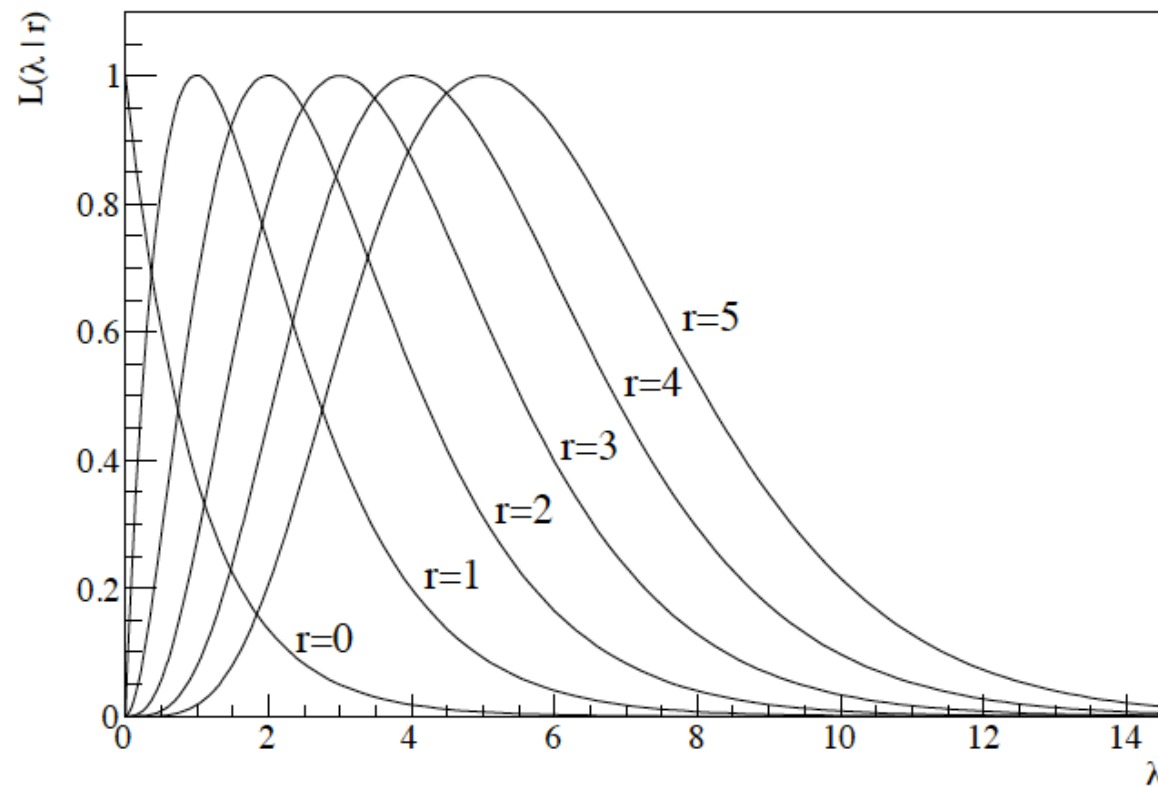
- ▶ For a given number of observed events resulting from the study of a rare process one wants to compute a limit on the true value of some underlying theory parameters (e.g. the mean occurrence of the rare process).
- ▶ This is a Poisson problem, where:

$$f(r, \lambda) = \frac{\lambda^r e^{-\lambda}}{r!},$$

$\lambda$       = mean/variance of the underlying distribution  
 $r$        = observed number of events



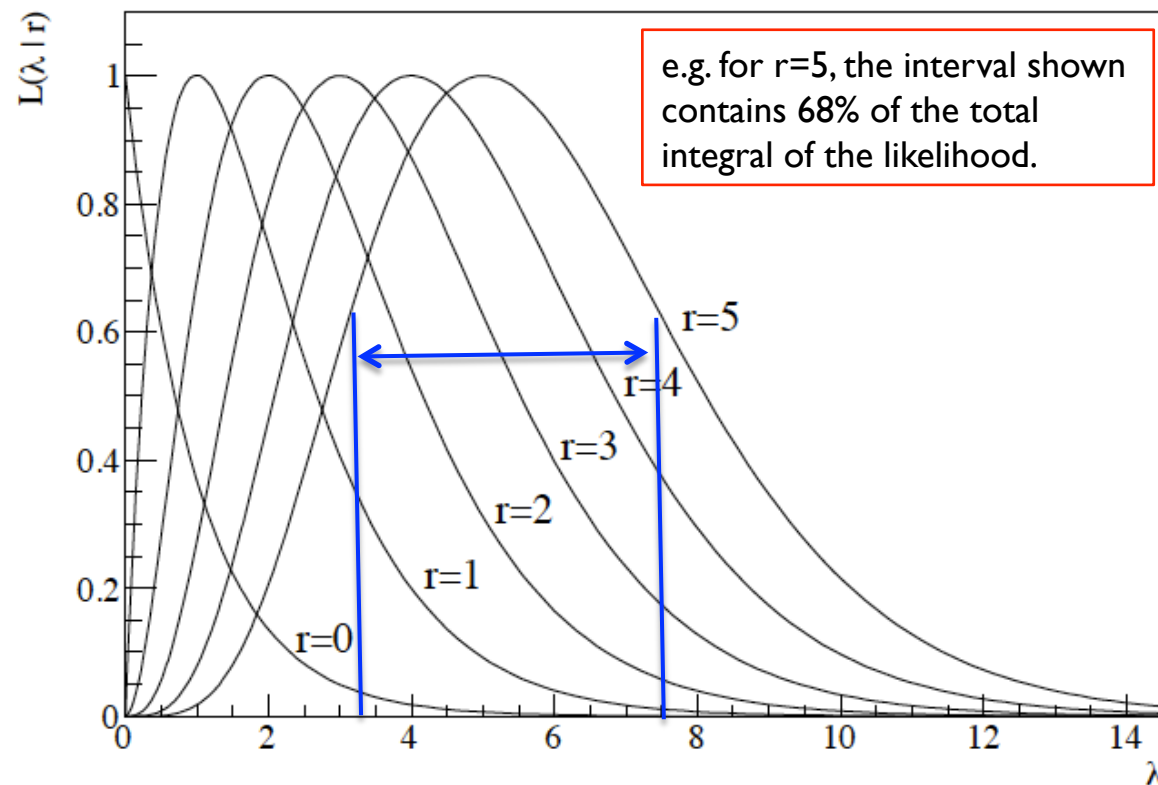
- For a given observation, we can compute a likelihood as a function of  $\lambda$ , e.g.



From each likelihood distribution one can construct a one or two sided interval by integrating  $L(\lambda, r)$  to obtain the desired coverage.



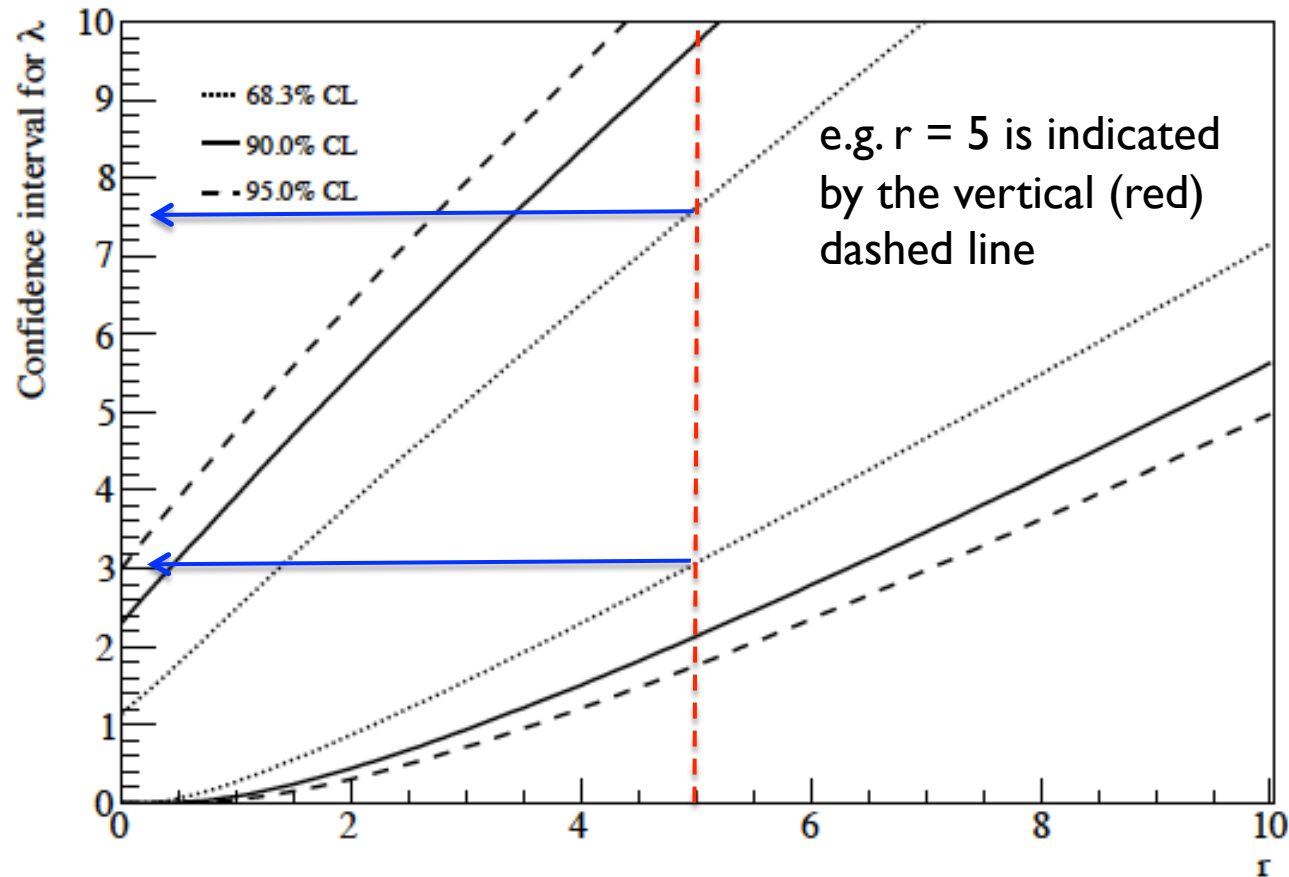
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From each likelihood distribution one can construct a one or two sided interval by integrating  $L(\lambda, r)$  to obtain the desired coverage.

We can build a 2D region from a family of likelihood curves for different values of  $r$ .





Physically there are discrete observed numbers of events for a background free process. If however background plays a role, then the problem becomes more complicated, and non-integer values may be of relevance..

# Multi-dimensional constraint

An example of constraining a model using data.

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email: [a.j.bevan@qmul.ac.uk](mailto:a.j.bevan@qmul.ac.uk)



# Motivation

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- ▶ The Standard Model of particle physics describes all phenomenon we know at a sub-atomic level.
- ▶ The model is incomplete:
  - ▶ Universal matter-antimatter asymmetry unknown
  - ▶ Nature of neutrinos unknown
  - ▶ What is Dark Matter
  - ▶ What is Dark Energy [related to a higher order GUT]
  - ▶ etc.
- ▶ Many particle physicists want to test the Standard Model precisely for two (related) reasons:
  - ▶ (i) understand the model better
  - ▶ (ii) see where it breaks...





# The problem

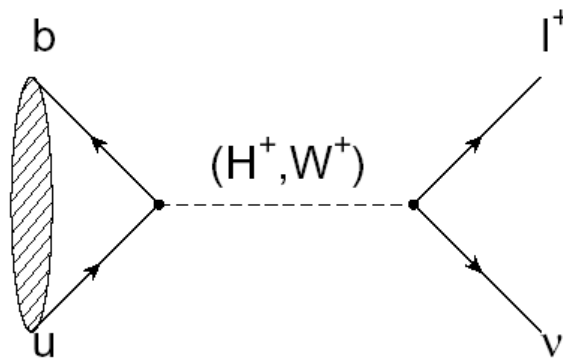
- ▶ The decay  $B^\pm \rightarrow \tau^\pm \nu$  has been measured, and can be compared with theoretical expectations.

- ▶ Measurement:

$$\mathcal{B}(B^\pm \rightarrow \tau^\pm \nu) = (1.15 \pm 0.23) \times 10^{-4}$$

- ▶ Standard Model expectation:

$$\mathcal{B}(B^\pm \rightarrow \tau^\pm \nu)_{SM} = (1.01 \pm 0.29) \times 10^{-4}$$



$$r_H = \frac{\mathcal{B}_{SM+NP}}{\mathcal{B}_{SM}}$$

For a simple extension of the Standard Model, called the type II 2 Higgs Doublet Model we know that  $r_H$  depends on the mass of a charged Higgs and another parameter,  $\beta$ .

$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$



# What can we learn about $m_H$ and $\tan\beta$ for this model?

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- ▶ We can compute  $r_H$  from our knowledge of the measured and predicted branching fractions:

$$r_H = 1.14 \pm 0.40$$

- ▶ How can we use this to constrain  $m_H$  and  $\tan\beta$ ?

$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$



## Method 1: $\chi^2$ approach

---

- Construct a  $\chi^2$  in terms of  $r_H$

$$\chi^2 = \left( \frac{r_H - \hat{r}_H(m_H, \tan \beta)}{\sigma_{r_H}} \right)^2$$





## Method 1: $\chi^2$ approach

- Construct a  $\chi^2$  in terms of  $r_H$

From SM theory  
and experimental  
measurement

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# Method 1: $\chi^2$ approach

- Construct a  $\chi^2$  in terms of  $r_H$

From SM theory  
and experimental  
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Calculate using

$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$

One has to select the parameter values.

$$\chi^2 = \left( \frac{r_H - \hat{r}_H(m_H, \tan \beta)}{\sigma_{r_H}} \right)^2$$

From SM theory  
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## Method 1: $\chi^2$ approach

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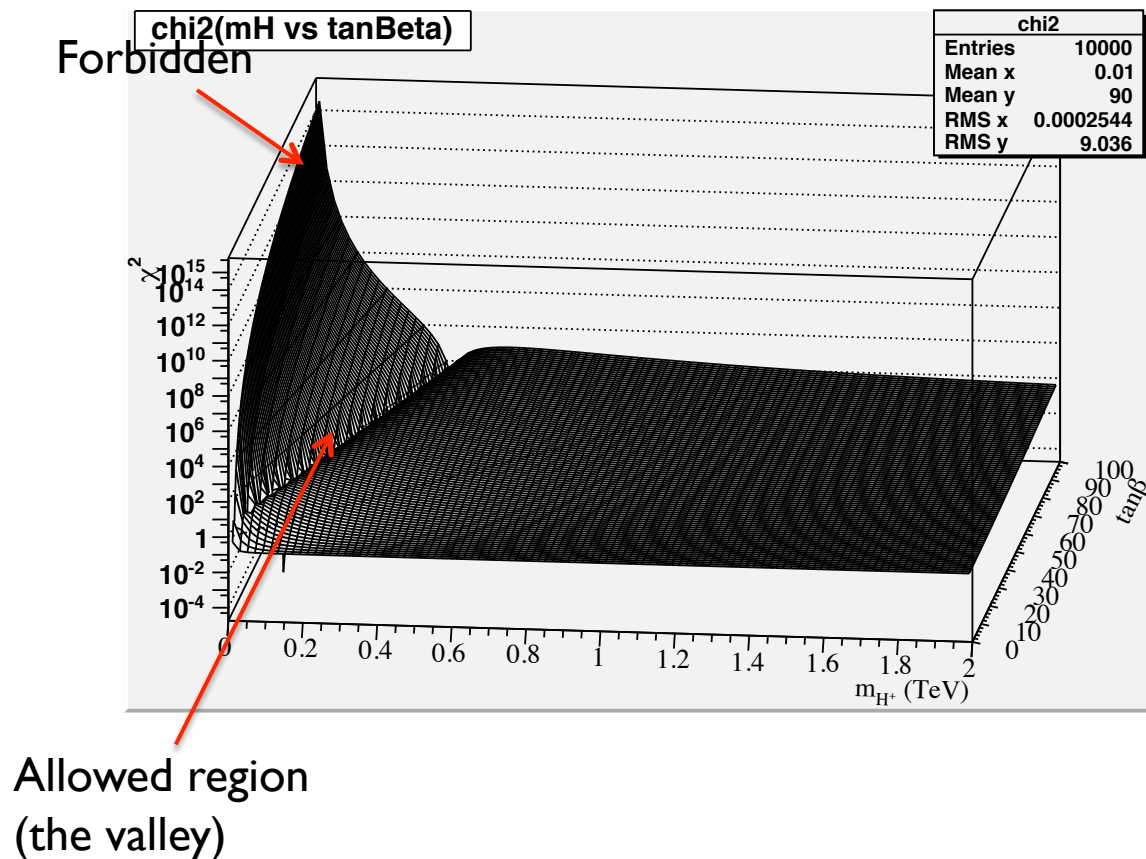
- ▶ For a given value of  $m_H$  and  $\tan\beta$  you can compute  $\chi^2$ .

- ▶ e.g.
$$\begin{aligned}m_H &= 0.2\text{TeV} \\ \tan\beta &= 10 \\ \hat{r}_H(m_H, \tan\beta) &= 0.93 \\ \chi^2 &= \left(\frac{1.14 - 0.93}{0.4}\right)^2 \\ &= 0.28\end{aligned}$$

- ▶ So the task at hand is to scan through values of the parameters in order to study the behaviour of constraint on  $r_H$ .



# Method 1: $\chi^2$ approach



A large  $\chi^2$  indicates a region of parameter space that is forbidden.

A small value is allowed.

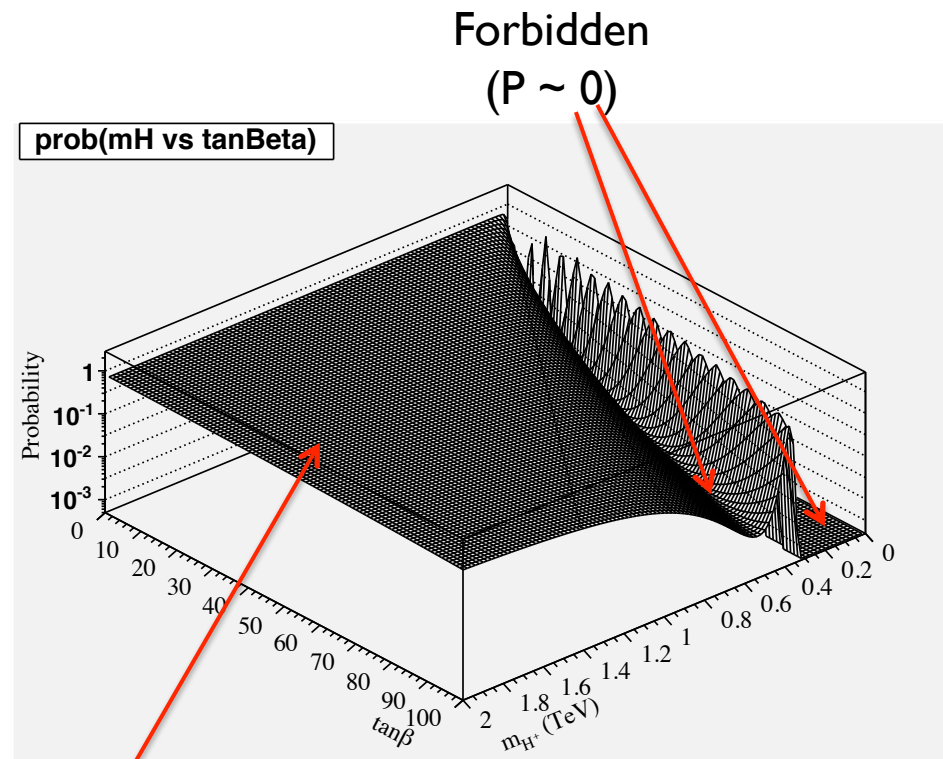
In between we have to decide on a confidence level that we use as a cut-off.

We really want to convert this distribution to a probability: so use the  $\chi^2$  probability distribution.

There are 2 parameters and one constraint (the data), so there are 2-1 degrees of freedom, i.e.  
 $\nu = 1$



# Method 1: $\chi^2$ approach



Allowed region  
( $P \sim 1$ )

Forbidden  
( $P \sim 0$ )

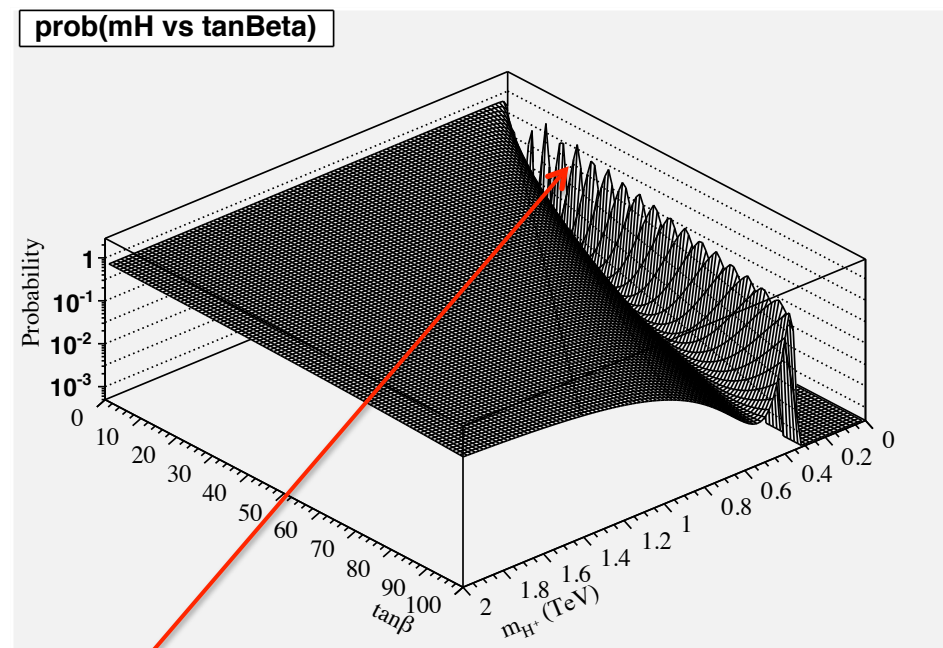
A  $P \sim 1$  means that we have no constraint on the value of the parameters (i.e. they are allowed).

A small value of  $P$ ,  $\sim 0$  means that there is a very low (or zero) probability of the parameters being able to take those values (i.e. the parameters are forbidden in that region).

Typically one sets a 1-CL corresponding to 1 or 3  $\sigma$  to talk about the uncertainty of a measurement, or indicate an exclusion region at that CL.



# Method 1: $\chi^2$ approach



Artefact: a remnant of binning the data. For these plots there are  $100 \times 100$  bins. As a result visual oddities can occur in regions where the probability (or  $\chi^2$ ) changes rapidly.

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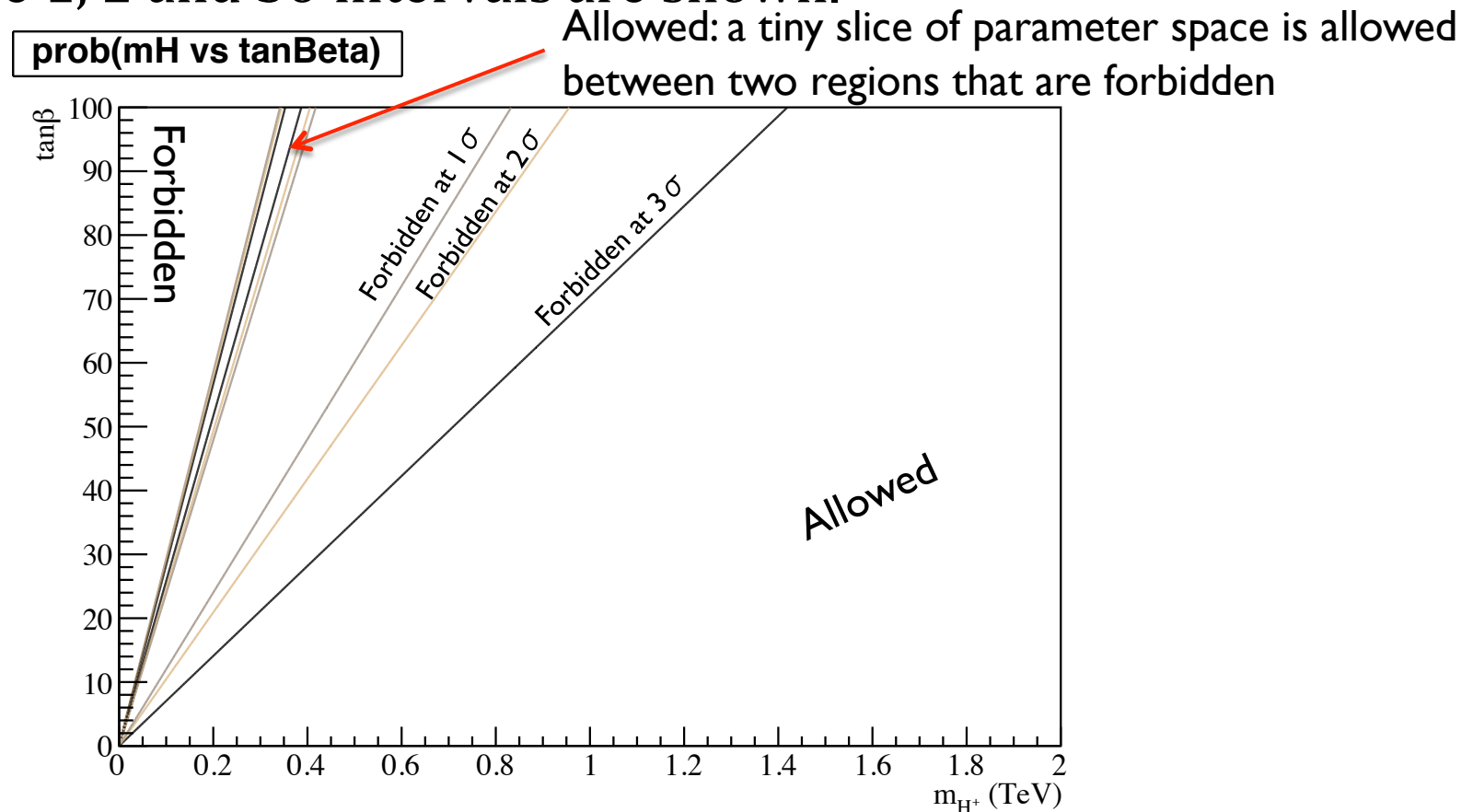
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# Method 1: $\chi^2$ approach

- ▶ A finer binning can be used to compute a 1-CL distribution. Here 1, 2 and 3 $\sigma$  intervals are shown.





## Summary: $\chi^2$ approach

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- ▶ This is nothing new – it is just a two-dimensional scan solution to a problem.
- ▶ It is however more computationally challenging to undertake (excel probably won't be a good fit to solving the problem):
  - ▶ 1D problem:  $N$  scan points
  - ▶ 2D problem:  $N^2$  scan points
  - ▶ As  $N$  becomes large (e.g. 100 or 1000) the number of sample points becomes very large.
    - ▶ i.e. the curse of dimensionality strikes.
- ▶ An MD problem has  $N^M$  sample points.
  - ▶ e.g. Minimal Super-Symmetric Model (MSSM) has  $\sim 160$  parameters, so one has a problem with  $N^{160}$  sample points. This is currently not a viable computational method to explore the complexities of MSSM.

