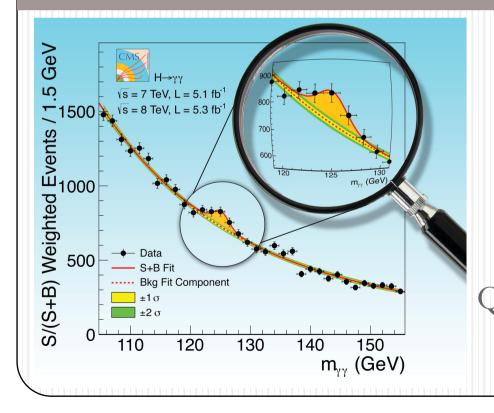


# Nature of the New Boson: a lonely Higgs, or one of many cousins?



Jim Olsen Princeton University

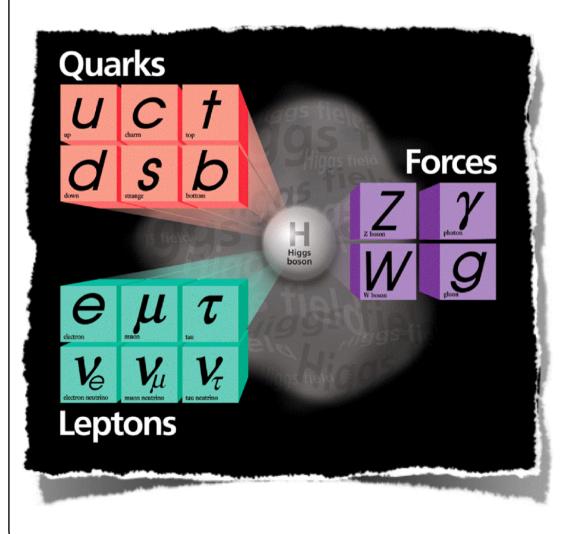
Colloquium of the School of Physics and Astronomy Queen Mary University of London January 17, 2014

### Outline

- Why do we need any Higgs bosons?
- Search for the Higgs boson
- Discovery of "a Higgs boson"
- Nature of the new boson
- Is there more than one Higgs boson?
- Future prospects in Higgs physics at the LHC

## Why do we need any Higgs bosons?

# One field to rule them all...



#### **Standard Model**

#### **Fundamental matter:**

quarks and leptons

**Symmetries:**  $U(1)_{\rm Y}$ ,  $SU(2)_{\rm L}$ ,  $SU(3)_{\rm C}$ 

Local gauge invariance: gauge bosons (force carriers)

#### Higgs field:

spontaneous EWK symmetry breaking and the Higgs boson

#### What's the problem?

Dirac Lagrangian:

"Local gauge invariance":  $\mathcal{L}$  should be invariant under  $\psi \rightarrow e^{i\theta(x)}\psi$ , but  $m_A^2 A^{\nu} A_{\nu}$  is not gauge invariant  $\rightarrow$  need  $m_A = 0$ !

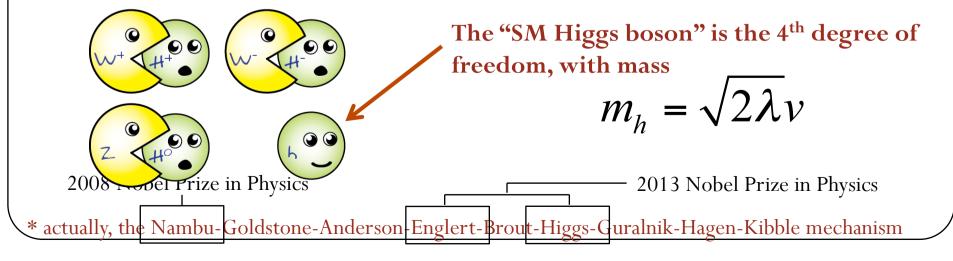
So it works for EM and QCD, but W and Z bosons are massive (~100 GeV). <u>Need a mechanism</u> to give mass to the W and Z.

The Higgs \* Mechanism (1964)  

$$\mathcal{L} = (D_{\mu}\phi)(D^{\mu}\phi) - (-\mu^{2}\phi^{2} + \lambda\phi^{4}) - \frac{1}{4}F^{\mu\nu}F_{\mu\nu}$$

$$\phi = \begin{pmatrix} \Phi_{1} \\ \Phi_{2} \end{pmatrix} \Rightarrow \begin{pmatrix} \nu + h \\ 0 \end{pmatrix}$$
Spontaneous symmetry  
breaking (v = "vev")
E Tando

3 degrees of freedom are "eaten" by the  $W^{\pm}$  and  $Z^{0}$ , thus acquiring mass, while the photon remains massless:



#### Properties of "the SM Higgs boson"

A single **elementary** scalar particle ( $J^P = 0^+$ ), quantum of the field that gives mass to the W and Z, plus quarks and leptons

Mass: 
$$m_h = \sqrt{2\lambda}v$$
  $v = \left(\sqrt{2}G_F\right)^{\frac{1}{2}} \approx 246 \,\text{GeV}$ 

Because  $\lambda$  is not predicted, the Higgs boson mass  $m_h$  is a free parameter in the Standard Model

Interaction: couples to particles according to their mass, including self-coupling (h\*  $\rightarrow$  hh)

$$m_W = \frac{gv}{2}, \ m_f = \frac{g_f v}{\sqrt{2}}$$

Is this the only possibility?

Additional Higgs fields 
$$\phi_i \longrightarrow 2m_W = \sum_i g_i v_i$$

Composite Higgs: top-quark condensate

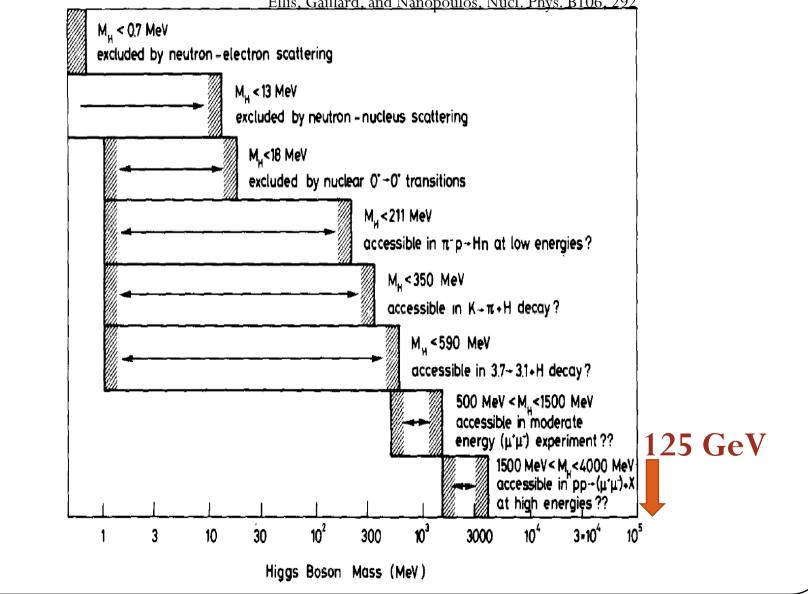
Technicolor: new gauge interactions

Extra dimensions

### Search for the Higgs boson

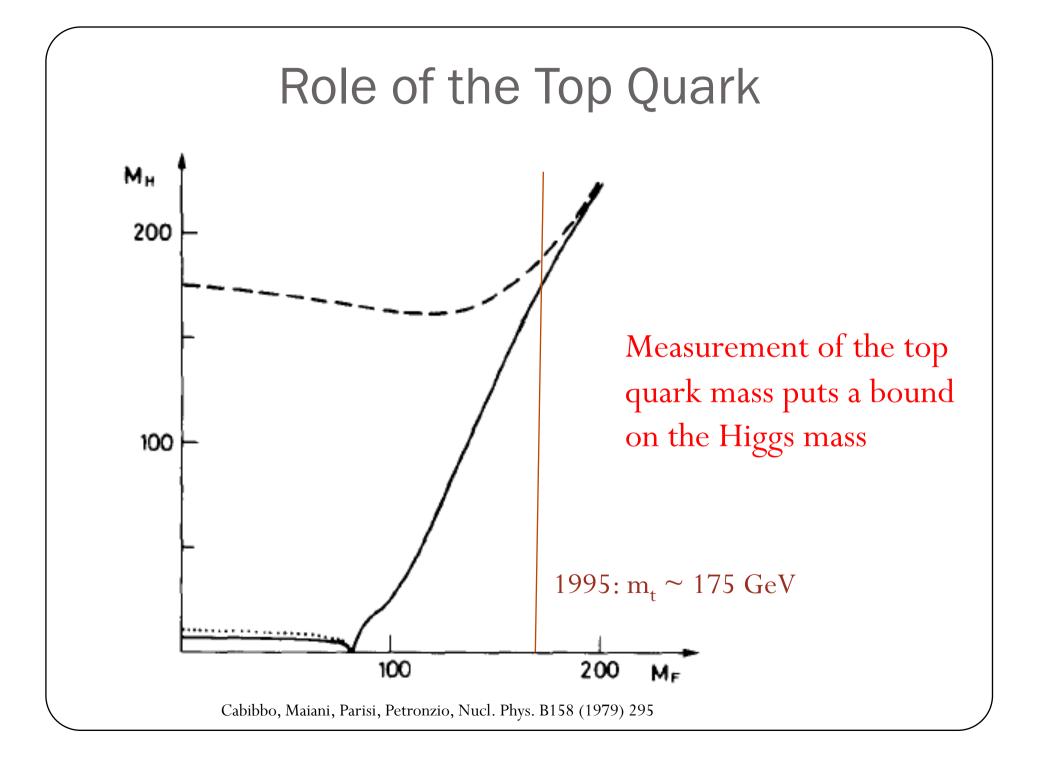
## Higgs Hunting: 1975

Ellis, Gaillard, and Nanopoulos, Nucl. Phys. B106, 292



#### Bounds on the SM Higgs Mass: ~1980 Mass of the Higgs Boson\* V(\$ Steven Weinberg Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138 (Received 15 December 1975) The stability of the vacuum sets a lower bound of order $\alpha G_F^{-1/2}$ on the Higgs-boson mass. For the simplest SU(2) $\otimes$ U(1) model, this lower bound is 1.738 $\alpha G_F^{-1/2}$ , or 3.72 GeV. and A. Linde, JETP Lett. 23 (1976) 64 Requiring V(v) < V(0) gives: $m_H > \left[\frac{3(2m_W^2 + m_Z^4) - 4\sum m_l^4 - 12\sum m_q^4}{16\pi^2}\right]^{\frac{1}{2}} > 7 \text{ GeV}$ PHYSICAL REVIEW D VOLUME 16, NUMBER 5 **1 SEPTEMBER 1977** Weak interactions at very high energies: The role of the Higgs-boson mass Benjamin W. Lee,\* C. Quigg,<sup>†</sup> and H. B. Thacker Fermi National Accelerator Laboratory, <sup>‡</sup> Batavia, Illinois 60510 (Received 20 April 1977) We give an S-matrix-theoretic demonstration that if the Higgs-boson mass exceeds $M_c = (8\pi\sqrt{2}/3G_F)^{1/2}$ , parital-wave unitarity is not respected by the tree diagrams for two-body scattering of gauge bosons, and the н weak interactions must become strong at high energies. We exhibit the relation of this bound to the structure of the Higgs-Goldstone Lagrangian, and speculate on the consequences of strongly coupled Higgs-Goldstone systems. Prospects for the observation of massive Higgs scalars are noted. $M_{\mu} \leq M_{c} = (8\pi\sqrt{2}/3G_{F})^{1/2} \simeq 1 \text{ TeV}/c^{2}$

By 1980 we knew the Higgs mass was  $7 \le m_h \le 1000 \text{ GeV}$ 



#### 1984: Higgs boson found, not!

#### Has the Higgs boson been seen in the Crystal Ball?

From the Crystal Ball detector at the DORIS electron-positron collider in Hamburg comes strong evidence of an 8.3-GeV particle for which there appears to be no prosaic explanation. "If it's real, it has to be very important," says Gordon Kane (University of Michigan), expressing the widespread excitement that this new state has generated among high-energy theorists.

The Crystal Ball group, an international collaboration of 13 laboratories. finds evidence for this uncharged particle, which they call the zeta, in the radiative decay of the upsilon meson, a bound state of the bottom quark (b) and its antiparticle (b) with a mass of 9.46 GeV (about 10 times the proton mass). If one didn't know better, the first explanation coming to mind would be that the zeta is simply the pseudoscalar ground state of the bb "bottomonium" system. But because the b quarks are so massive (about 5 GeV), theorists can apply nonrelativistic potential models to calculate the mass spectrum of bb bound states with considerable confidence. Such calculations clearly predict that the bb ground state should lie only about 40 MeV below the upsilon (9.46 GeV). The zeta (8.32 GeV) is more than a full GeV too light. With similiar confidence in calculations of the charmonium spectrum, they dismiss the possibility that the zeta could be a highlying cc state.

Higgs scalars. What most excites the high-energy community is the very real possibility that the zeta may be the long-sought-after Higgs particle. The standard theory of elementary particles makes almost no detailed predictions about the properties of the Higgs particle except to insist that it (or perhaps even several) must appear somewhere in the mass range between a few GeV and a few TeV.

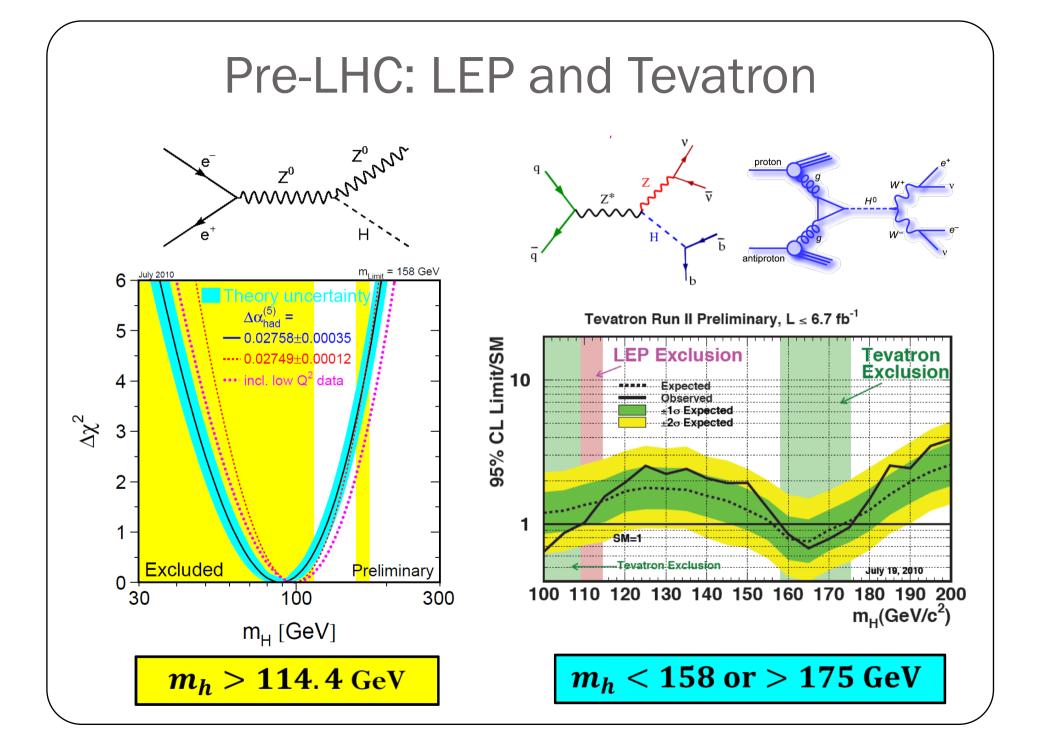
The predicted Higgs particle is a manifestation of the spontaneous-symmetry-breaking mechanism that permits the electromagnetic and weak interactions to look so different from one another despite the underlying symmetry of the Glashow-Salam-Weinberg electroweak theory that unifies them. This mechanism, invoking scalar fields with nonvanishing vacuum expectation values, was introduced twenty years ago by Peter Higgs (University of Edinburgh) to deal with the finite mass of the exchanged gauge boson in an early attempt at a gauge field theory of the strong interactions. We now know that any renormalizable gauge theory must have such additional Higgs fields if its gauge field quanta (the photon and the vector bosons W ± and Z<sup>0</sup>, for example) and its fermions (the quarks and leptons) are not all to be massless.

The minimal Higgs mechanism for the electroweak theory predicts just one

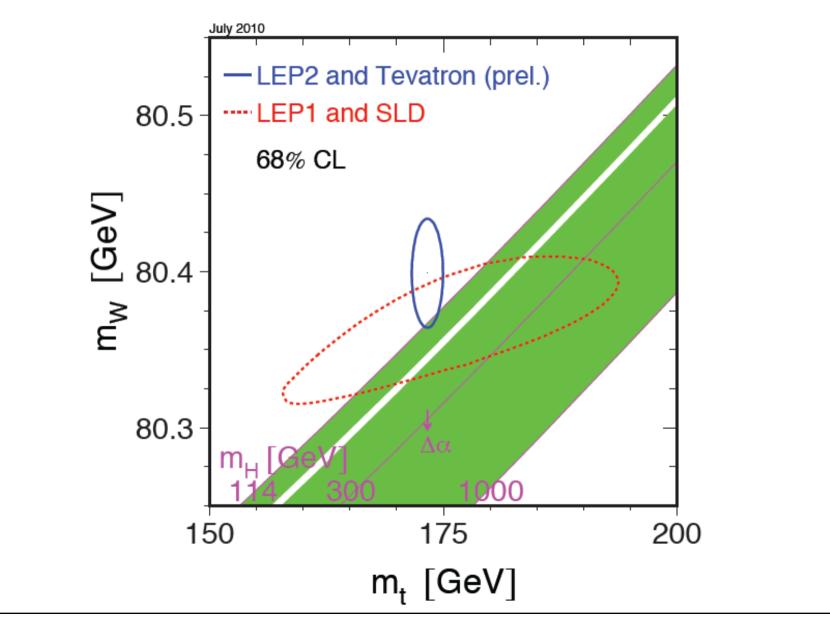
plex scalar Higgs field four independent Higgs fields. When nents vanish, one ha metrical state in v bosons W± and Zº weak interaction are the photon, and all di the weak and electro tions disappears. Bu state is not the low um" state. In the t the Higgs fields choose nonzero valferromagnet choose ground state despit symmetry of its Har This spontaneous metry breaking is i occurred as a phase the cooling of the ea the Higgs field sough gy, its four compone nonzero vacuum ex three of them becom nal components of bosons, giving the W large masses (80-95 ( terize the weak inter ton remains massless netic field is purely remaining fourth His came the Higgs par experimenters are lo What would this H

doublet (in weak isos



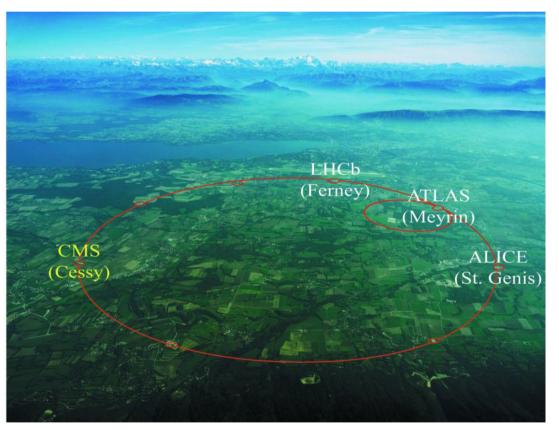


#### At the dawn of the LHC era



# Discovery of "a Higgs boson"

#### Large Hadron Collider



pp collider inside the 27km LEP tunnel (1998-2008)

Highest energy collider ever built (8 trillion electron volts)

1232 superconducting dipole magnets with B > 8 Tesla

Magnets are colder than space

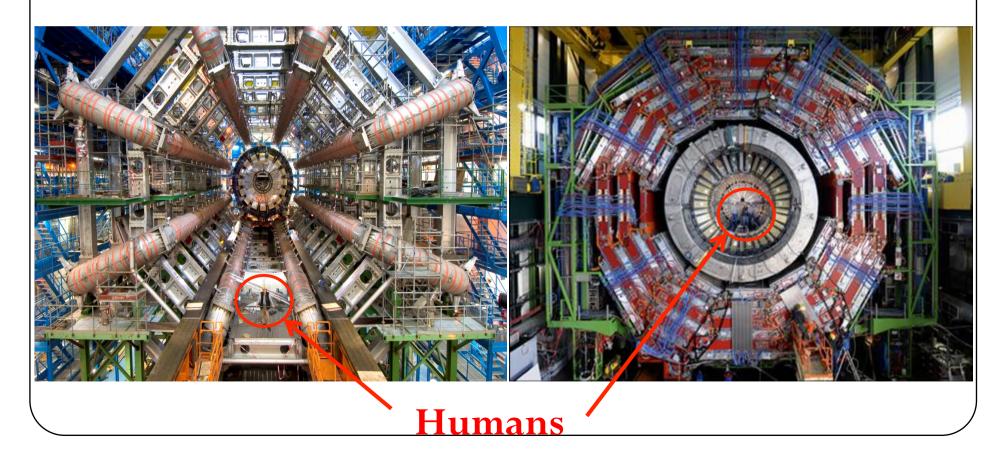
World's largest cryogenic plant

Roughly 750,000 Higgs bosons have been produced by the LHC

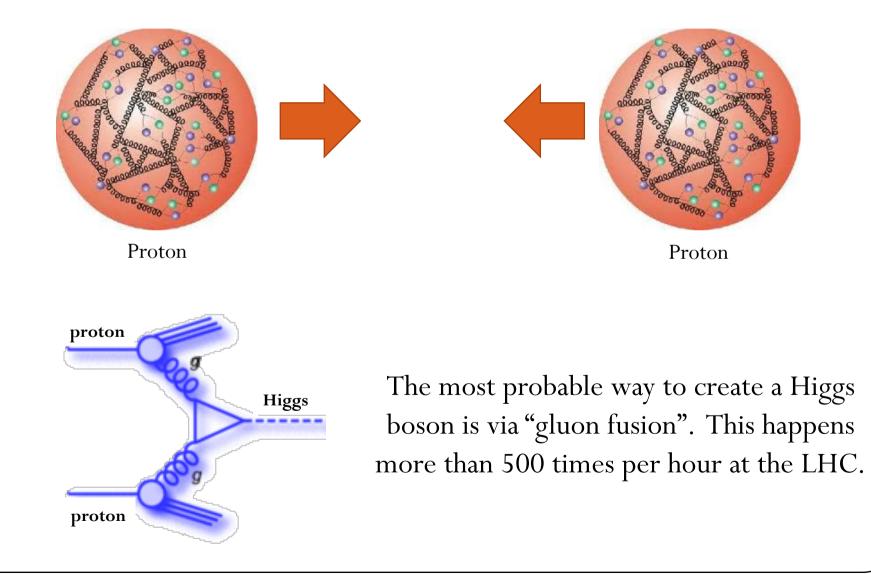
#### ATLAS and CMS

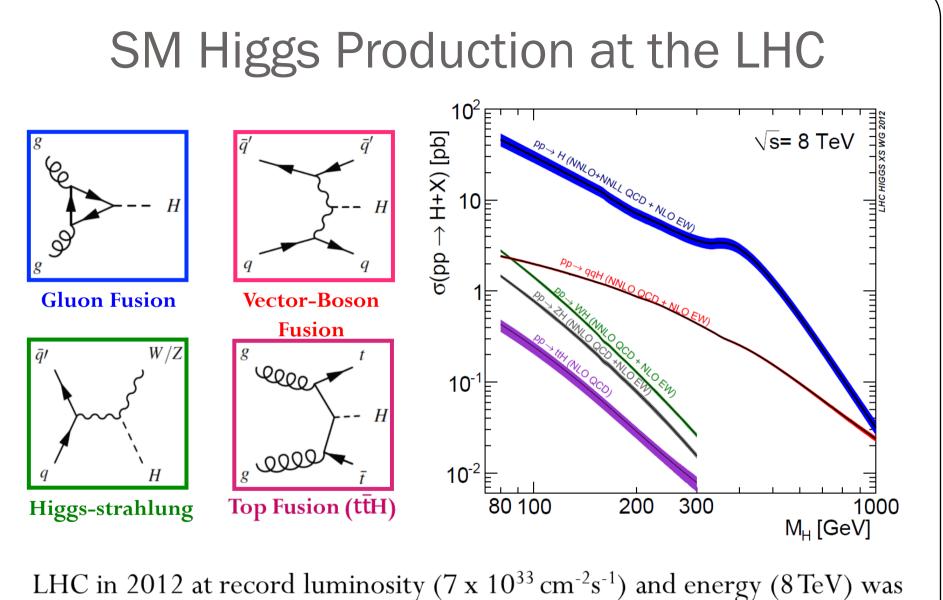
~3000 scientist, engineers, and students working on **each** experiment

Giant multipurpose particle detectors designed to find or exclude the Higgs boson and signs of physics beyond the SM

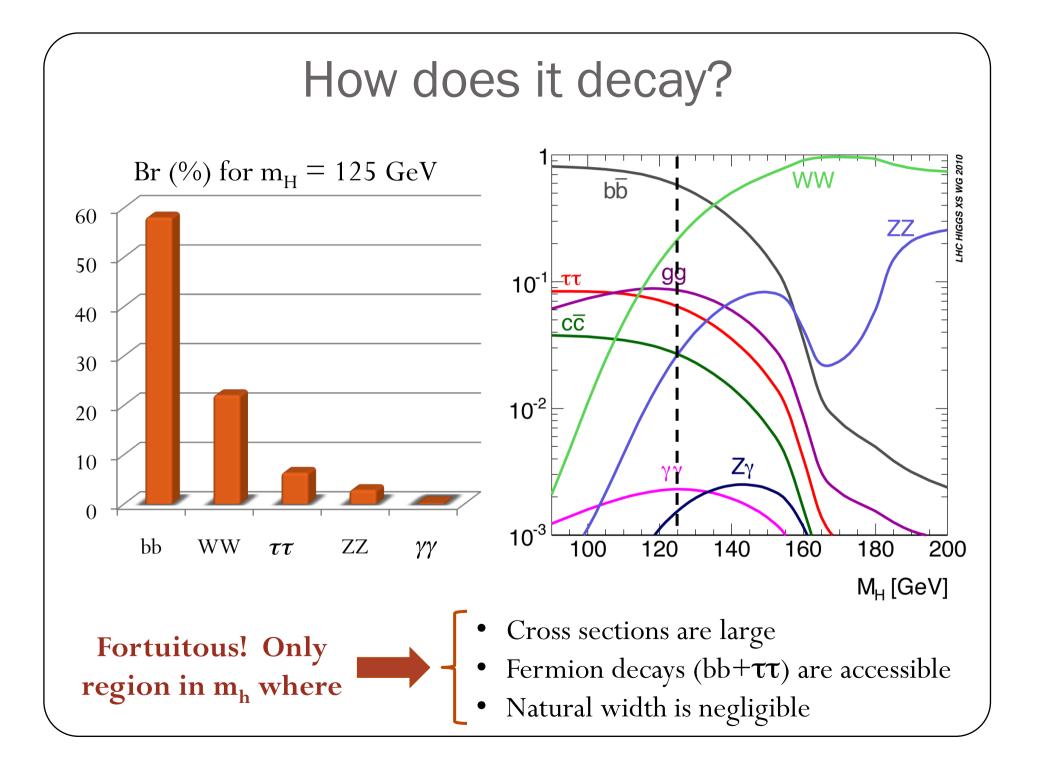


## Creating a Higgs boson





producing SM Higgs bosons ( $M_H = 125$  GeV) at a rate ~750/hr



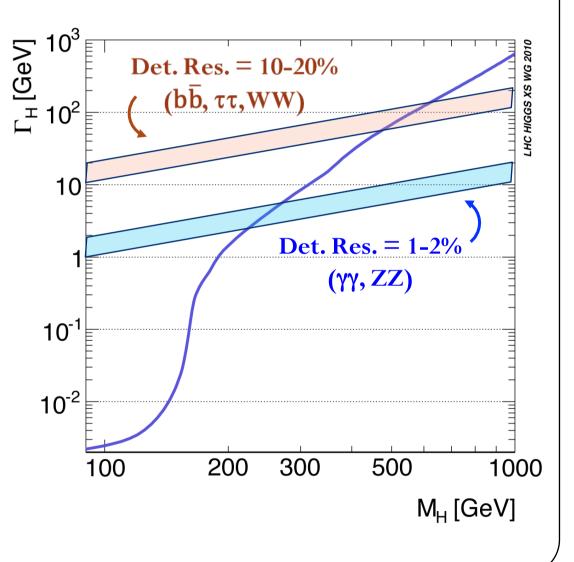
#### What does it look like?

(a) Low mass Narrow!  $\Gamma_{\rm H}/m_{\rm H} \sim 10^{-5}$ Observed width dominated by *detector resolution* 

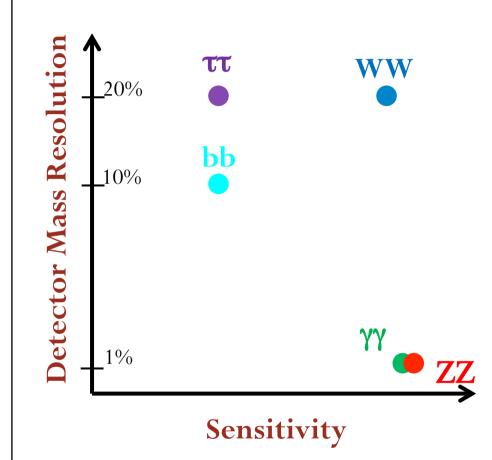
*a***High mass** 

Higgs becomes a broad resonance dominated by *natural width* 

Theory input is critical



#### SM Higgs Analyses @ LHC



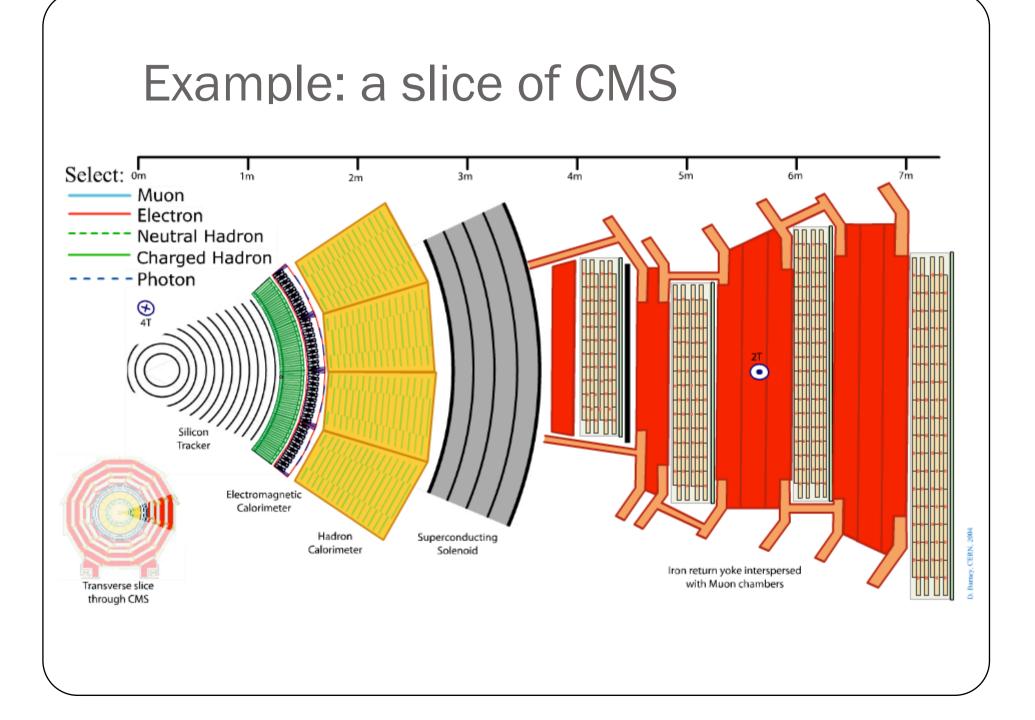
Three classes of SM analysis:

1) High sensitivity, high resolution  $\gamma\gamma$ , ZZ

2) High sensitivity, low resolutionww

3) Low sensitivity, low resolutionττ, bb

Need **multi-purpose detectors** like ATLAS and CMS to find the Higgs boson at the LHC!



# pile-up (PU)

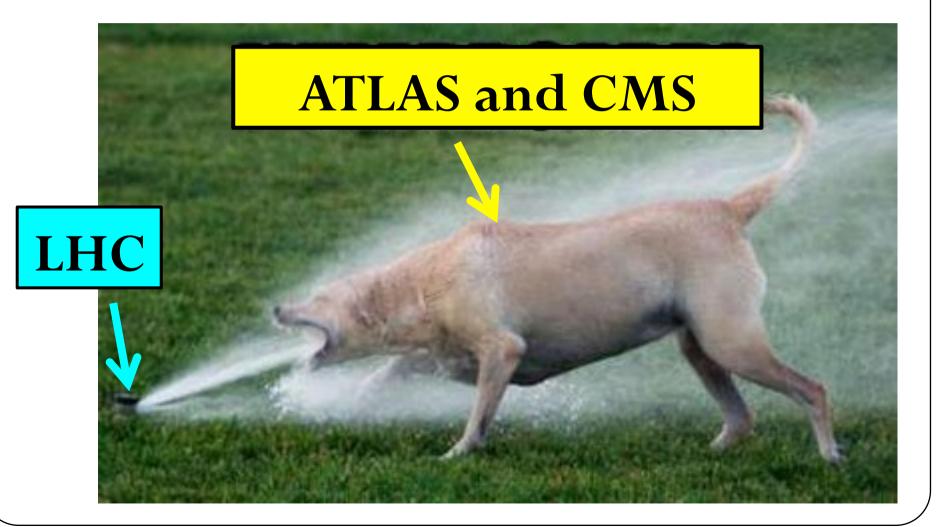
At this high luminosity, multiple collisions per beam-crossing occur.

Experimental challenge to cope with high PU.

Reconstruction and analyses are designed to be robust against PU.

78 reconstructed vertices

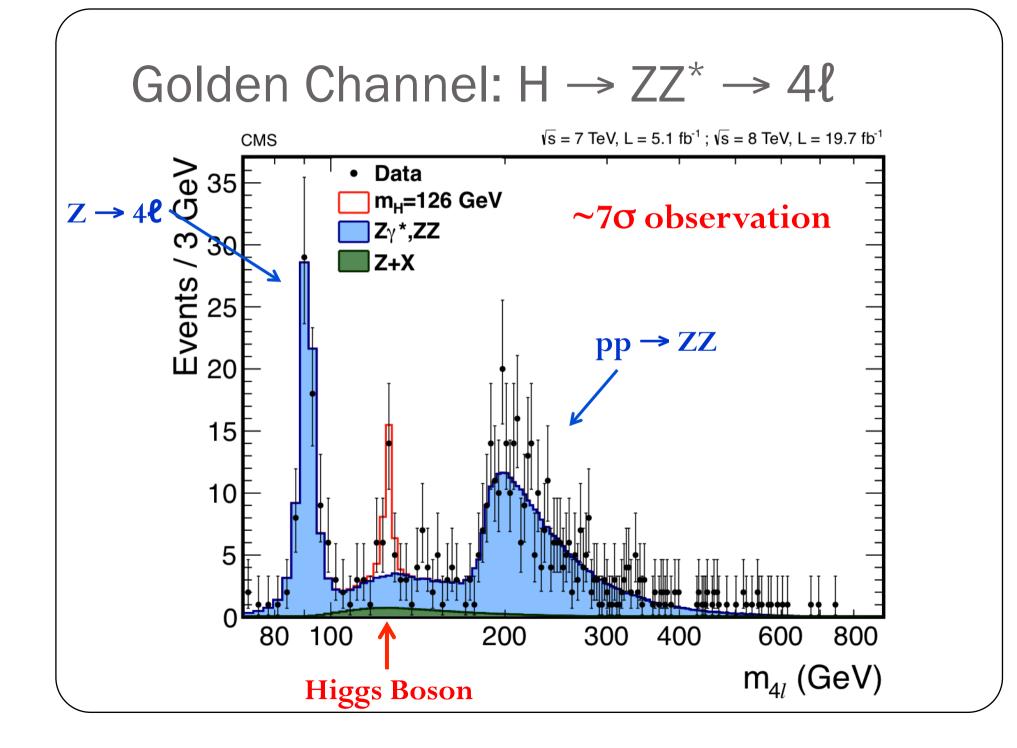
What it's like to be an experimentalist at the LHC:

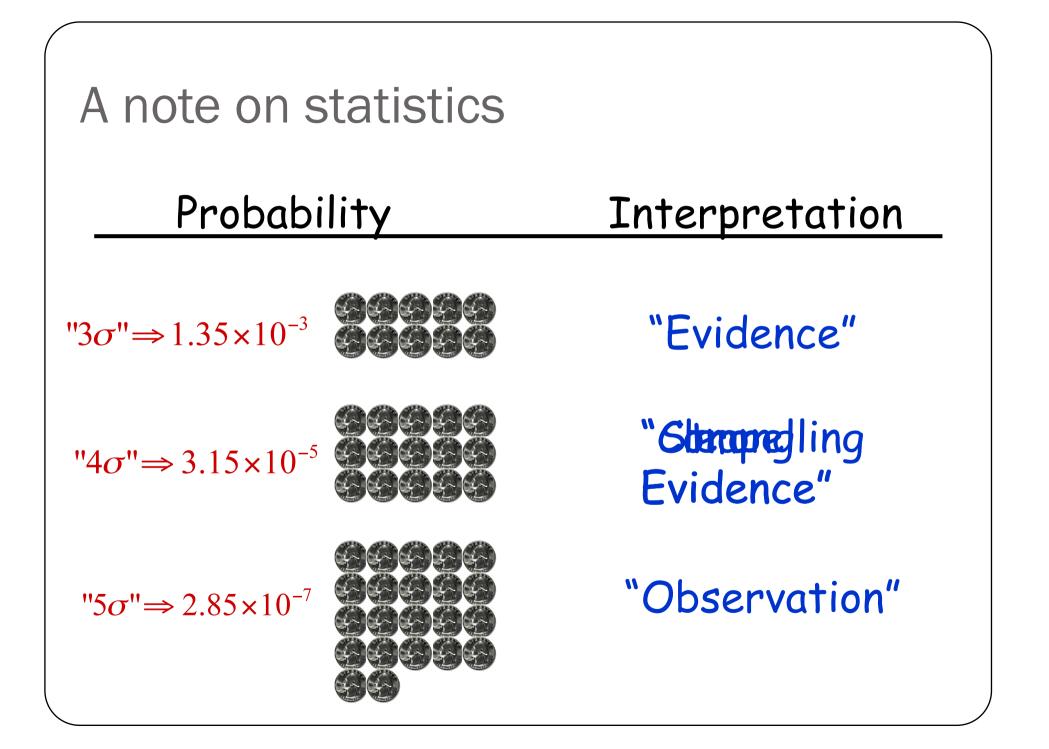




CMS Experiment at LHC, CERN Data recorded: Thu Oct 13 03:39:46 2011 CEST Run/Event: 178421 / 87514902 Lumi section: 86







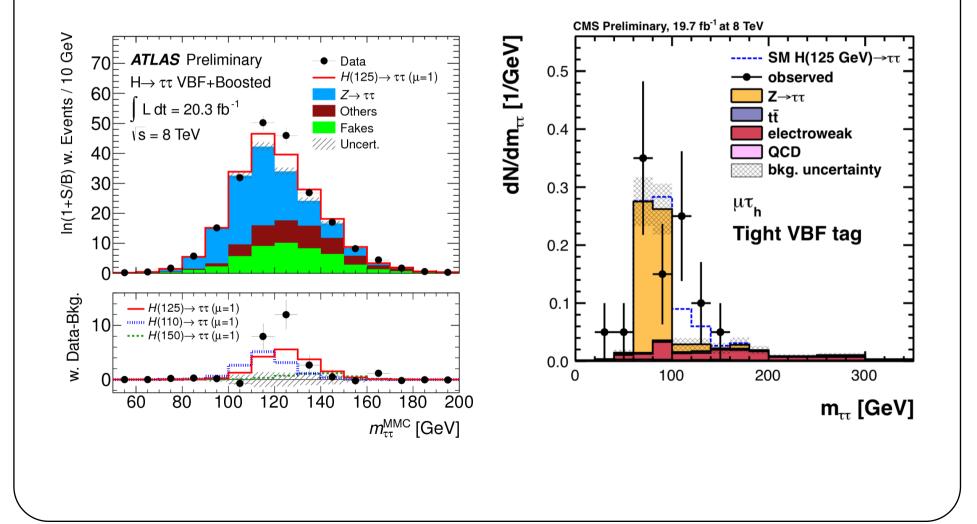
#### Significance ( $\sigma$ ) by decay channel

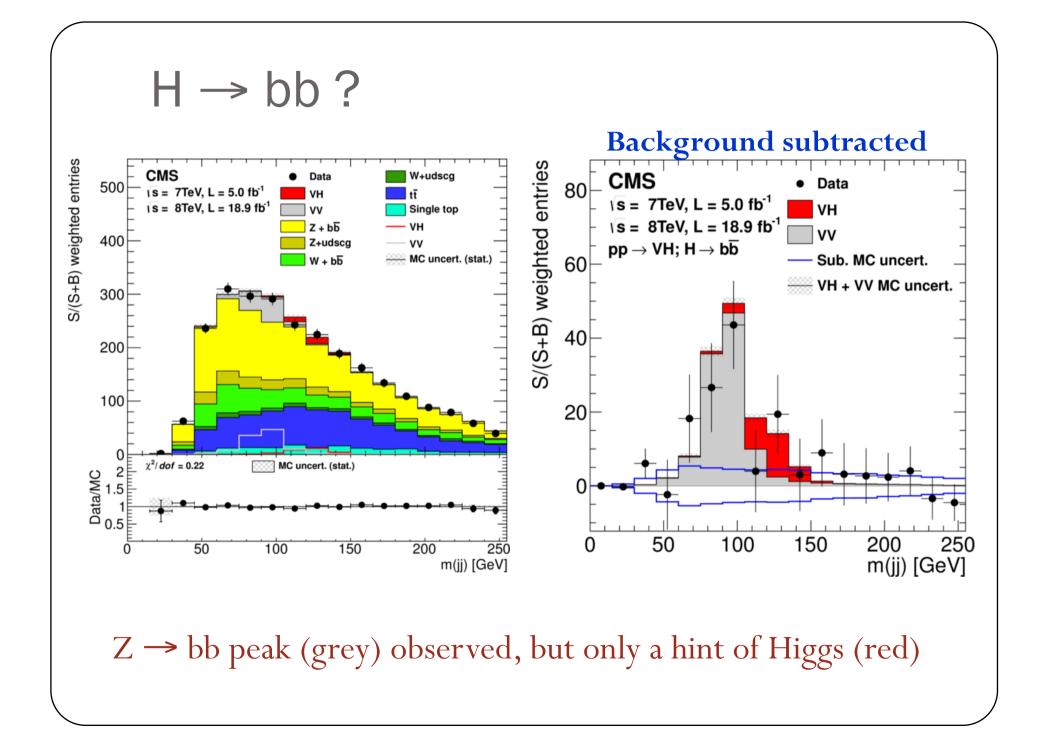
Channel	ATLAS (expected)	ATLAS (observed)	CMS (expected)	CMS (observed)
h → γγ	4.1	7.4	4.2	3.2
$h \rightarrow ZZ$	4.4	6.6	6.7	6.8
h→WW	3.7	3.8	5.8	4.3
$h \rightarrow \tau \tau$	3.2	4.1	3.6	3.4
$h \rightarrow bb$	1.6	~0	2.1	2.1

First "clear, strong, compelling" evidence for fermionic decays

#### $H \rightarrow \tau \tau$ : ATLAS and CMS

End of 2013: ATLAS and CMS showed updated results on the full LHC dataset



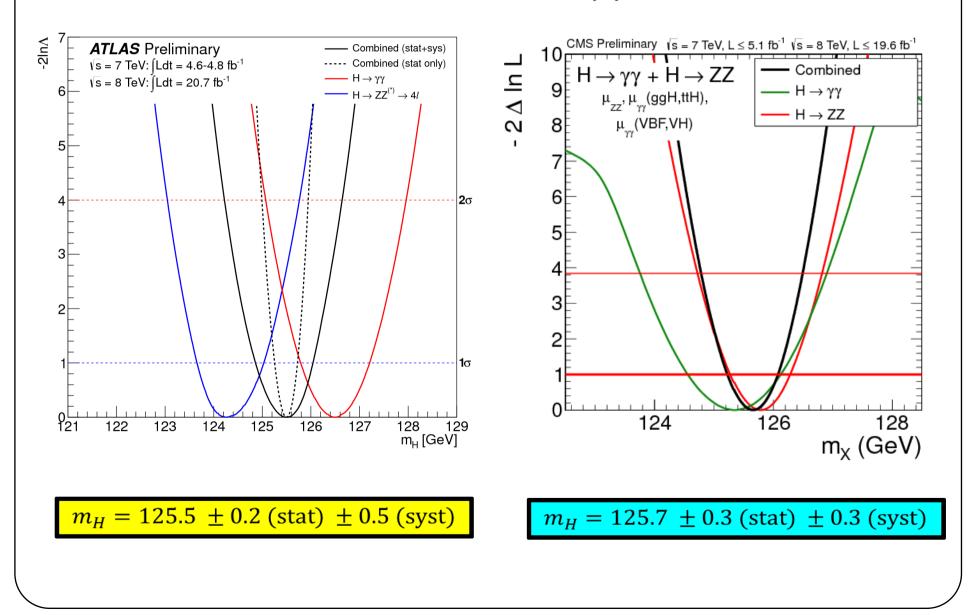


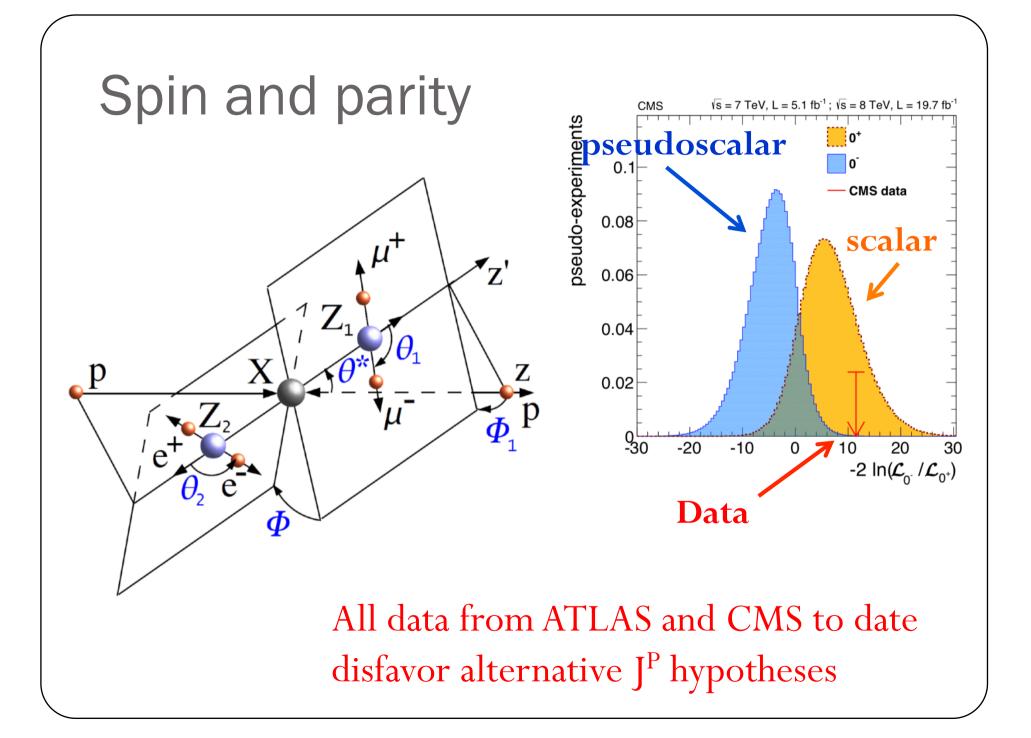
#### Nature of the New Boson

#### Signal strength by decay channel SM $\mu \equiv \sigma \times \text{Br} / \sigma_{\text{SM}} \times \text{Br}_{\text{SM}}$ $\sqrt{s} = 7 \text{ TeV}, L \le 5.1 \text{ fb}^{-1} \sqrt{s} = 8 \text{ TeV}, L \le 19.6 \text{ fb}^{-1}$ т<sub>н</sub> = 125.7 GeV CMS Preliminary Combined **ATLAS** Preliminary m<sub>u</sub> = 125.5 GeV $\mu = 0.80 \pm 0.14$ р<sub>SM</sub> = 0.65 $\mu = 0.80 \pm 0.14$ W.Z H $\rightarrow$ bb $\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.7 \text{ fb}^{-1}$ $H \rightarrow bb$ $\sqrt{s} = 8 \text{ TeV}$ : $\int Ldt = 13 \text{ fb}^{-1}$ $\mu = 1.15 \pm 0.62$ $H \rightarrow \tau \tau$ $\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.6 \text{ fb}^{-1}$ $v_{s} = 8 \text{ TeV}: \int Ldt = 13 \text{ fb}^{-1}$ $H \rightarrow WW^{(*)} \rightarrow hh$ $H \rightarrow \tau \tau$ Vs = 7 TeV: ∫Ldt = 4.6 fb<sup>-1</sup> $\mu = 1.10 \pm 0.41$ $\sqrt{s} = 8 \text{ TeV}$ : $\int Ldt = 20.7 \text{ fb}^{-1}$ $H \rightarrow \gamma \gamma$ $H \rightarrow \gamma \gamma$ $V_{s} = 7 \text{ TeV}: \int Ldt = 4.8 \text{ fb}^{-1}$ $\sqrt{s} = 8 \text{ TeV}$ : $\int Ldt = 20.7 \text{ fb}^{-1}$ $\mu = 0.77 \pm 0.27$ $H \rightarrow ZZ^{(*)} \rightarrow 4$ √s = 7 TeV: ∫Ldt = 4.6 fb<sup>-1</sup> $\sqrt{s} = 8 \text{ TeV}: \int Ldt = 20.7 \text{ fb}^{-1}$ $H \rightarrow WW$ $\mu = 0.68 \pm 0.20$ $\mu = 1.30 \pm 0.20$ Combined $\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.6 - 4.8 \text{ fb}^{-1}$ $\sqrt{s} = 8 \text{ TeV};$ $\int Ldt = 13 - 20.7 \text{ fb}^{-1}$ $H \rightarrow ZZ$ $\mu = 0.92 \pm 0.28$ -1 0 +1 2.5 0.5 1.5 2 0 Signal strength $(\mu)$ Best fit σ/σ<sub>sm</sub>

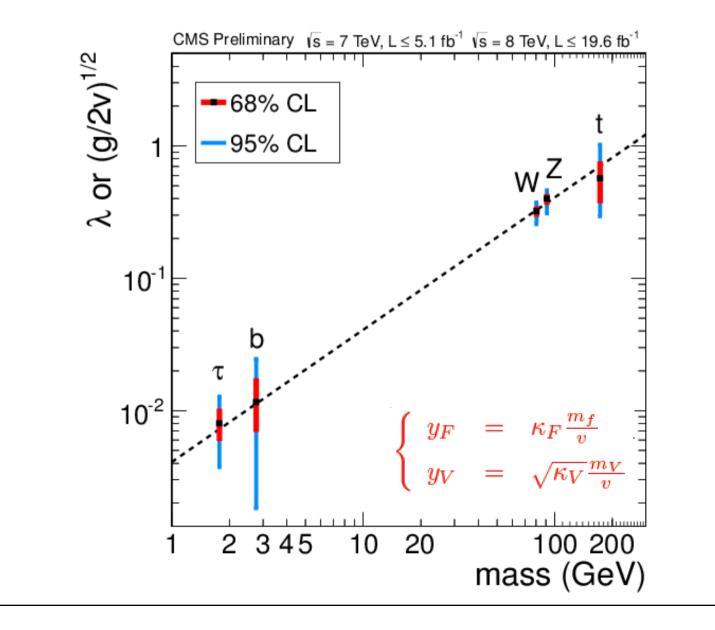
Both experiments are consistent with the SM expectation

#### Mass (GeV) from $\gamma\gamma$ and ZZ





## Couplings of the Higgs boson



#### New results indicate that particle discovered at CERN is a Higgs boson

ISSUES AND EXPERTS

- CERN

#### Higgs boson and new pope confirmed

March 14, 2013

💟 Tweet 🛛 🛐 Facebook 🛛 🖂 Email 🖱 Print

#### 'God particle' is for real

Scientists are confirming that a new subatomic particle discovered at the world's most powerful particle accelerator is indeed an elusive Higgs boson, also referred to as 'the God particle.' It was discovered during experiments at the Large Hadron Collider (LHC) at CERN, Switzerland last July. Scientists, who say they have a "long way to go" to know what kind of Higgs boson it is, are reporting the confirmation at the Moriond physics conference in Italy this week. The Higgs boson is the only particle in the Standard Model of Physics that has never been observed. SFU physicists **Dugan O'Neil, Bernd Stelzer** and **Michel Vetterli** are involved with ATLAS – one of two international physics experiments involving the LHC– and can comment on the news. Vetterli is currently at CERN where he'll spend the next year and can also do Skype interviews.

Michel Vetterli, +41 22 767 4368; vetz@triumf.ca; mikevetterli (Skype) Dugan O'Neil, 778.782.5623; dugan\_oneil@sfu.ca Bernd Stelzer, 778.782.7731; stelzer@sfu.ca Background: http://at.sfu.ca/eUnDFV

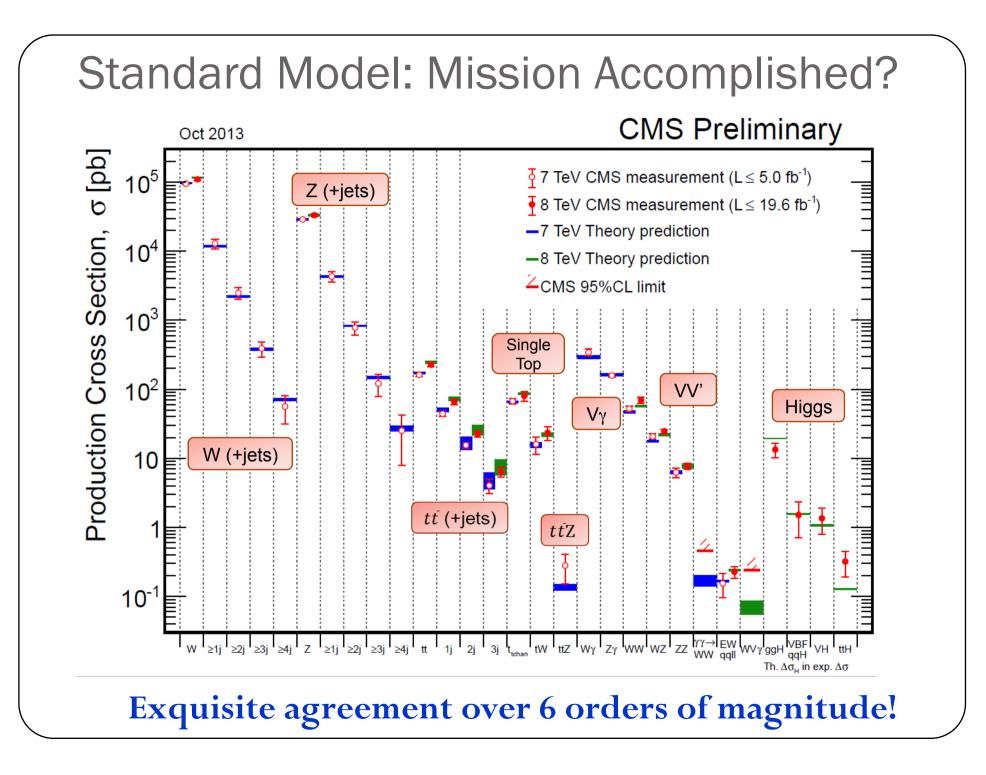
#### First non-European pope chosen

Jorge Mario Bergoglio of Argentina was elected yesterday as the new pope to lead the Roman Catholic Church. **Donald Grayston**, a retired religious studies professor at SFU, says that it is interesting the new pope comes from the new world.

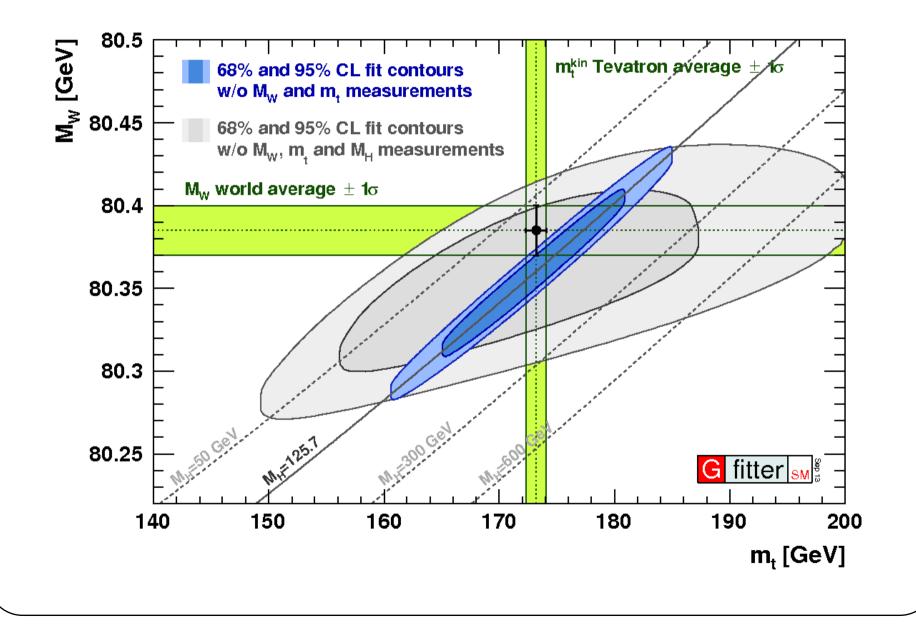
"I don't think the Curia is ready to let go of its stranglehold on change," says Grayston. "But I'm betting that the non-Curia cardinals wanted to elect someone who could tackle the Curia. I see it as very appropriate, given that 40-percent of the world's Catholics live in Latin America."



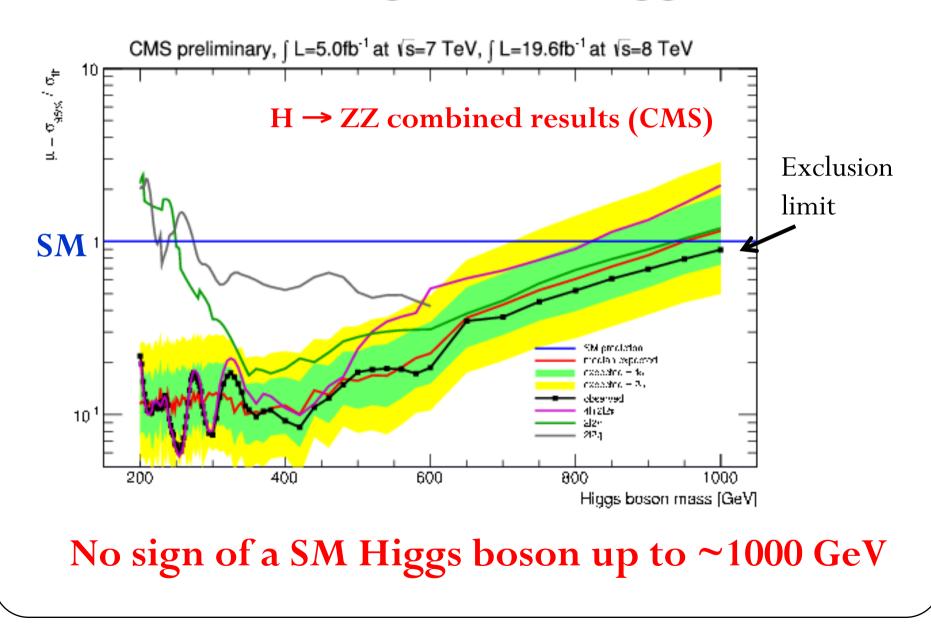
# Is there more than one Higgs boson? (Is there **anything** else besides SM?)



### Standard Model: Mission Accomplished?



## Search for a high-mass Higgs Boson



### Search for "everything else"

#### Supersymmetry

 $\underline{\mathbf{e}, \mu, \tau, \gamma}$  Jets  $\mathsf{E}_{\mathsf{T}}^{\text{miss}} \int \mathcal{L} dt [fb^{-1}]$ 

Yes Yes Yes Yes Yes

Yes Yes Yes Yes Yes Yes

Yes Yes Yes Yes Yes Yes Yes 19 Yes Yes Yes 20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.7 20.7

> Yes Yes Yes Yes Yes Yes 20.3 20.3 20.7 20.7 20.7 20.7 20.3

Yes

Yes

4.6 4.6 4.7 20.7 20.7

20.3

2-6 jets 3-6 jets 7-10 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 0-3 jets 0-2 jets

3 b 7-10 jets 3 b 3 b Yes Yes Yes Yes 20.1 20.3 20.1 20.1

2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b

0

2.6

1 jet 1-5 jets Yes Yes 20.3 22.9 15.9 4.7 20.3

7 jets Yes Yes Yes

6-7 jets

2 e, µ (SS) 0-3 b

Mass limit

275-430 GeV

m Gov

600 GoV

ATLAS SUSY Searches\* - 95% CL Lower Limits

0 1 e,µ

0 1 e,μ 2 e,μ 2 e,μ 1-2 τ 2 γ 1 e,μ + γ

0-1 e,μ 0-1 e,μ

0 2 e, µ (SS) 1-2 e, µ 2 e, µ 2 e, µ 0 1 e, µ 0

2 e, μ (Z) 3 e, μ (Z) 1 5

2 e,μ 2 e,μ

2τ 3 e,μ 3 e,μ 1 e,μ

1 a displ vb

γ 1 b 2 e, μ (Z) 0-3 jets 0 mono-jet

Model

MSUGRA/CMSS MSUGRA/CMSS MSUGRA/CMSS

Gravitino LSP

 $\tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ 

 $\tilde{t}_{L,R}\tilde{t}_{L,R}, \tilde{t} \rightarrow t\tilde{x}_1^0$   $\tilde{t}_1\tilde{x}_1, \tilde{x}_1 \rightarrow t\tilde{v}(t\tilde{v})$   $\tilde{t}_1\tilde{x}_1, \tilde{x}_1 \rightarrow t\tilde{v}(t\tilde{v})$ 

 $\tilde{x} \rightarrow \tilde{h} t, \tilde{h} \rightarrow hs$ 

 $\gamma \tilde{l}_{L} \gamma \tilde{l}_{L} \ell (\tilde{\gamma} \gamma), \ell \tilde{\gamma} \tilde{l}_{L} \ell (\tilde{\gamma} \gamma)$   $\gamma W \tilde{\gamma}^{0} Z \tilde{\gamma}^{0}$ 

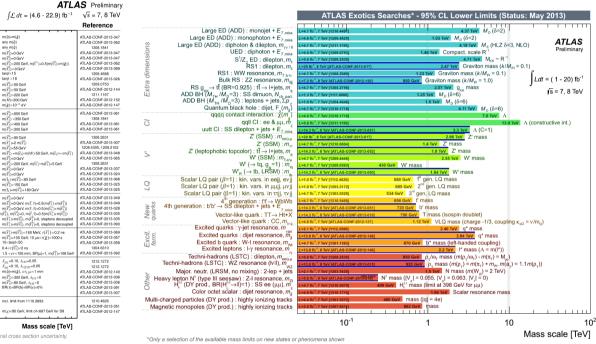
Direct  $\tilde{\chi}_1^+ \tilde{\chi}_1^-$  prod., long-lived  $\tilde{\chi}_1^\pm$  Disapp. trk Stable, stopped  $\tilde{g}$  R-hadron 0

 $\begin{array}{l} \frac{\mathrm{N}q, \epsilon \to \tau_{\mathrm{e}}}{\mathrm{LFV}\,pp \to \widetilde{v}_{\tau} + X, \widetilde{v}_{\tau} \to e + \mu} & 2 \, e, \mu \\ \mathrm{LFV}\,pp \to \widetilde{v}_{\tau} + X, \widetilde{v}_{\tau} \to e(\mu) + \tau & 1 \, e, \mu + \tau \\ \mathrm{Binear}\,\mathrm{RPV}\,\mathrm{CMSSM} & 1 \, e, \mu \\ \widetilde{x}_{1}^{+}\widetilde{x}_{1}^{-}, \widetilde{x}_{1}^{+} \to W \widetilde{x}_{1}^{0}, \widetilde{x}_{1}^{0} \to ee\widetilde{v}_{\mu}, e\mu\widetilde{v}_{e} & 4 \, e, \mu \end{array}$ 

 $\begin{array}{c} & \text{ar RPV CMSSM} \\ & \tilde{\kappa}_1^+ \rightarrow W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e e \tilde{y}_{\mu}, e \mu \tilde{y}_e & 4 \cdot e, \mu \\ & \tilde{\chi}_1^+ \rightarrow W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau r \tilde{y}_e, e \tau \tilde{y}_\tau & 3 \cdot e, \mu + \tau \\ & 0 \end{array}$ 

 $\ddot{g} \rightarrow b \bar{b} \tilde{\chi}_{1}^{0}$   $\ddot{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$   $\ddot{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$   $\ddot{g} \rightarrow b \bar{t} \tilde{\chi}_{1}^{+}$ 

 $\begin{array}{l} \mathsf{MSUGAWCMSSM}\\ \mathsf{q}, \; \overline{q} \rightarrow q^{T}_{1}\\ \overline{g}, \; \overline{g} \rightarrow q^{T}_{1}\\ \overline{g}, \; \overline{g} \rightarrow q^{T}_{1}\\ \overline{g}, \; \overline{g} \rightarrow qq^{T}_{1} \rightarrow qqW^{+}p^{0}_{1}\\ \overline{g}, \; \overline{g}, \; \overline{g} \rightarrow qqW^{+}(\mathcal{H}(\mathcal{H}(\mathcal{H})\mathcal{H})^{T}_{1}\\ \mathsf{GMSB}(\mathcal{H}(\mathcal{LSP})\\ \mathsf{GMSB}(\mathcal{H}(\mathcal{LSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP}))\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP}))\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP})\\ \mathsf{GGM}(\mathsf{bign}(\mathcal{HSP}))\\ \mathsf{GSM}(\mathsf{bign}(\mathcal{HSP}))\\ \mathsf$ 



"Exotica"



n(#)>200 Ge n(ĝ)>10<sup>-4</sup> eV

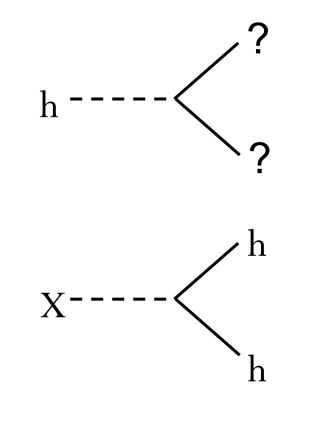
 $m(\tilde{t}_{1}^{0}) < 600 \text{ GeV}$   $m(\tilde{t}_{1}^{0}) < 350 \text{ GeV}$   $m(\tilde{t}_{1}^{0}) < 400 \text{ GeV}$   $m(\tilde{t}_{1}^{0}) < 300 \text{ GeV}$ 

 $\begin{array}{l} m(\tilde{t}_{1}^{*}) < 90 \; \text{GeV} \\ m(\tilde{t}_{1}^{*}) = 2 \; m(\tilde{t}_{1}^{*}) \\ m(\tilde{t}_{1}^{*}) = 55 \; \text{GeV} \\ m(\tilde{t}_{1}^{*}) = m(\tilde{t}_{1}) - r \\ m(\tilde{t}_{1}^{*}) = 0 \; \text{GeV} \end{array}$ 

Many many many searches, no sign of any physics beyond the SM

**Experimental Perspective** 

We have a new particle, h(126), use it as a 'scout' to search for other particles:



Exotic (invisible) decays, e.g. dark matter

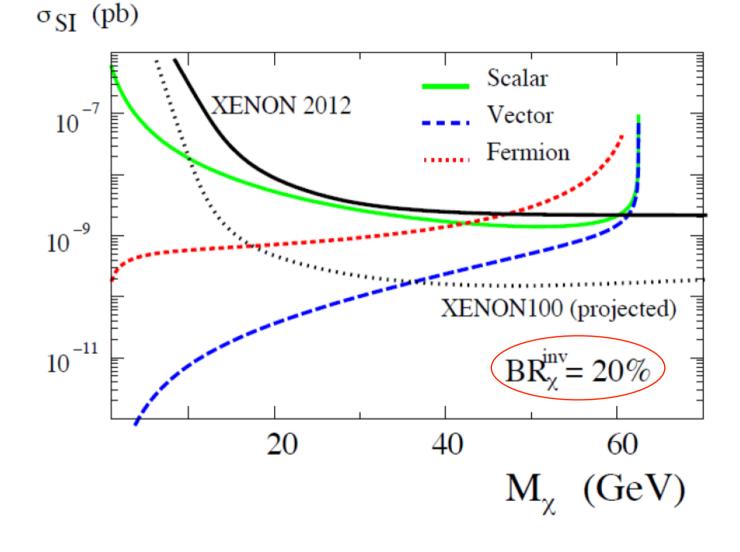
Heavy resonances (Higgs, SUSY, Exotica) decaying to the observed Higgs boson

# **Theoretical Perspective**

- The SM is not complete, there are specific questions to answer. Some examples:
  - Higgs mass: is our universe 'fine-tuned'?
  - Why is the EWK scale so far from the GUT scale?
  - What is the nature of dark matter?
- An extended Higgs sector could help in answering some of these questions
  - Models with additional Higgs bosons:
    - MSSM: H, h, A,  $H^+$ ,  $H^-$
    - NMSSM: add a singlet
  - "Higgs portal" coupling to dark matter

ATLAS and CMS are searching for all of these

# Dark matter search: $h(126) \rightarrow invisible$



Best limit from the LHC so far (CMS): BR<sub>inv</sub> < 0.54 (95% C.L.)

# Additional Higgs bosons

Extensions of the SM (e.g., SUSY) naturally include additional scalar doublets, leading to multiple Higgs bosons:

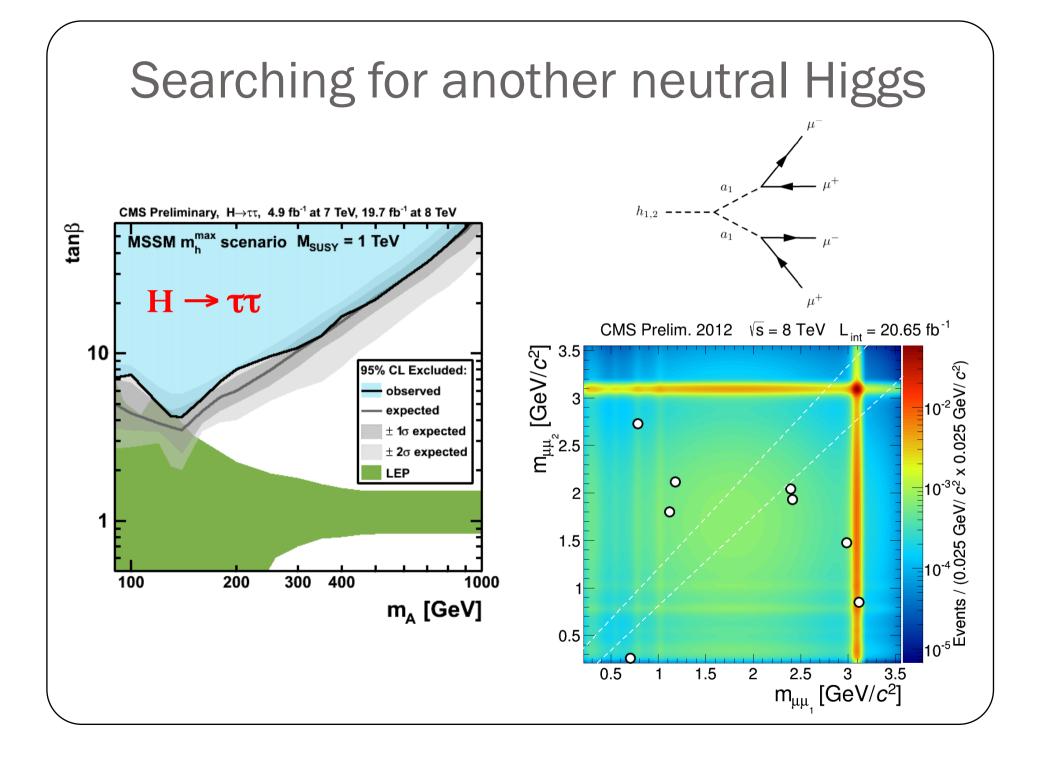
 $\phi_1, \phi_2$ : 8 degrees of freedom



three go to the W, Z masses, leaving five physics Higgs bosons

#### H<sup>+</sup>, H<sup>-</sup>, H, h, A

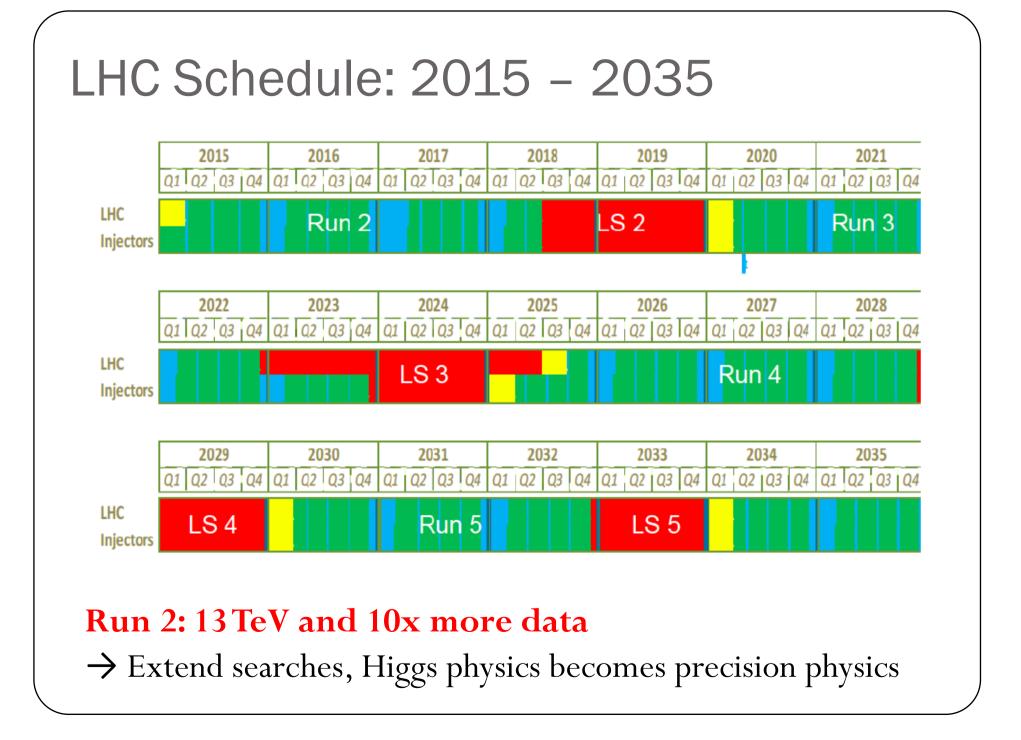
The Higgs boson we found could be H or h (looks like A is ruled out from spin-parity measurements), we search for additional Higgs bosons above and below h(126).



## Summary

- The new particle @ "126 GeV" is observed to decay to all gauge bosons, and in the right proportion
- Consistent mass between CMS and ATLAS
- Now "clear, strong, compelling" evidence for  $H \rightarrow \tau \tau$
- Spin-parity measurements disfavor alternative hypotheses
- Signal strength and couplings consistent with the SM
- No sign for any other SM-like Higgs boson
- No sign of (any of) the beyond-SM Higgs bosons

If it is not the SM Higgs boson it certainly is a good actor! Continuing studies of h(126), looking for new Higgs particles





"It is not the end. It is not even the beginning of the end. It is, perhaps, the end of the beginning." –Winston Churchill