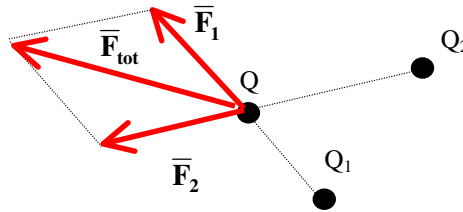


Answer 1

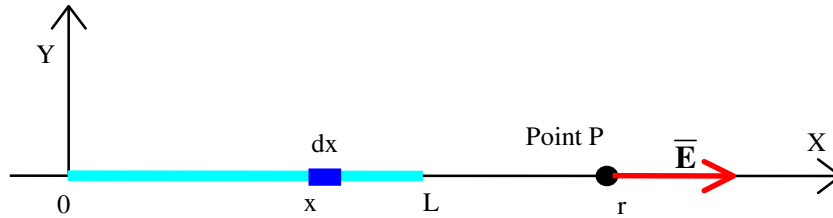
- (a) Principle of Superposition: The force on a charge due to a number of other charges is given by the vector sum of all the forces on it due to the individual charges.

e.g.: Total force on Q is

$$\vec{F}_{\text{tot}} = \vec{F}_1 + \vec{F}_2$$



(b)



Clearly, \vec{E} points along the +x direction (the direction in which a positive charge would tend to move if placed at the point).

To find the magnitude of \vec{E} , consider a small element of the rod, dx , at a distance x from the origin. We can regard this as a point charge, find its contribution to the field at the point in question, and then integrate over the whole rod to find the total field.

The charge on dx is $dQ = \lambda dx$.

The contribution of dx to E at point P is

$$dE = \frac{\lambda dx}{4\pi\epsilon_0 (r-x)^2}$$

So the total field at P is

$$E = \left[\frac{\lambda}{4\pi\epsilon_0} \right] \int_0^L \frac{1}{(r-x)^2} dx$$

Let $u = r - x$ $du = -dx$ $x = 0, u = r$ $x = L, u = r - L$

Therefore

$$E = - \left[\frac{\lambda}{4\pi\epsilon_0} \right] \int_r^{r-L} \frac{1}{u^2} du = - \left[\frac{\lambda}{4\pi\epsilon_0} \right] \left[-\frac{1}{u} \right]_r^{r-L} = \left[\frac{\lambda}{4\pi\epsilon_0} \right] \left[\frac{1}{u} \right]_r^{r-L}$$

So

$$E = \frac{\lambda L}{4\pi\epsilon_0 r(r-L)}.$$

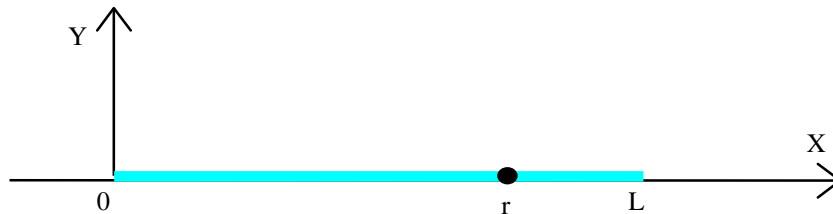
As noted above, \vec{E} points in the x-direction, so we write

$$\vec{E} = \frac{\lambda L}{4\pi\epsilon_0 r(r-L)} \hat{i}.$$

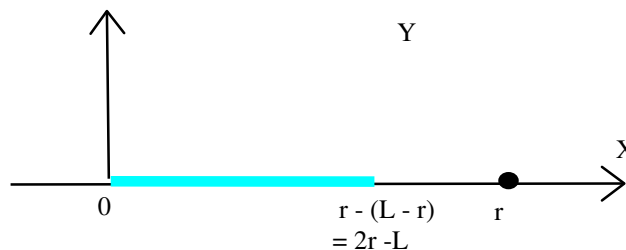
If $r \gg L$ then $\vec{E} \approx \frac{\lambda L}{4\pi\epsilon_0 r^2} \hat{i} \approx \frac{Q}{4\pi\epsilon_0 r^2} \hat{i},$

where Q is the total charge on the rod. This is consistent with the fact that the rod appears as a point charge for $r \gg L$, so that the electric field is given by Coulmb's Law.

- (c) Consider a point inside the rod, a distance r from the origin. By symmetry, the contributions to the electric field from the sections of length $L - r$ on either side of it will cancel out.



So this is equivalent to:



We can now use the above result with L replaced by $2r - L$:

$$\vec{E} = \frac{\lambda(2r - L)}{4\pi\epsilon_0 r(r - (2r - L))} \hat{i} = \frac{\lambda(2r - L)}{4\pi\epsilon_0 r(L - r)} \hat{i}.$$

Answer 2

(a) Gauss' s Law: $\Phi = \oint \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$

where Φ = Electric flux through an arbitrary closed surface

\vec{E} = Electric field at the surface

$d\vec{A}$ = Normal vector to an infinitesimal area, dA , on the surface

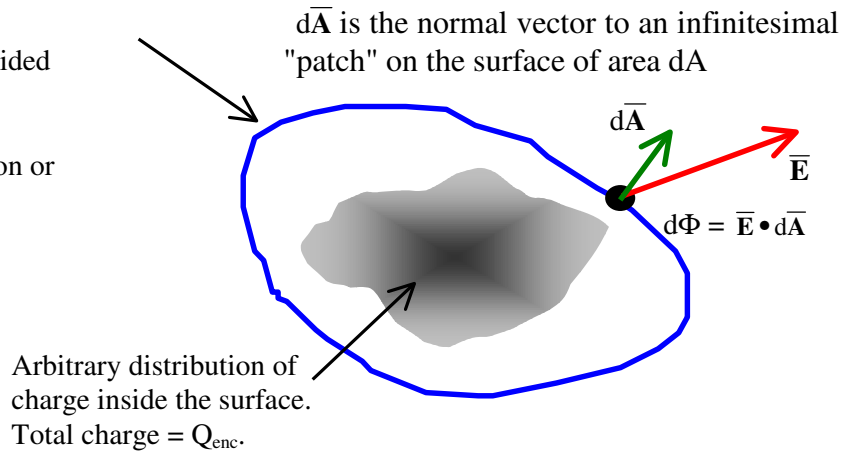
Q_{enc} = Total electric charge contained within the closed surface

ϵ_0 = Permittivity constant

$\oint \Rightarrow$ Integral over a *closed* surface

Gauss' s law states that the total electric flux, Φ , through a **closed surface** is equal to the enclosed charge divided by ϵ_0 .

This holds regardless of the location or distribution of the charge inside.



(b)

1. Field pattern: By symmetry, the electric field lines due to a spherically symmetric charge distribution must be radial.

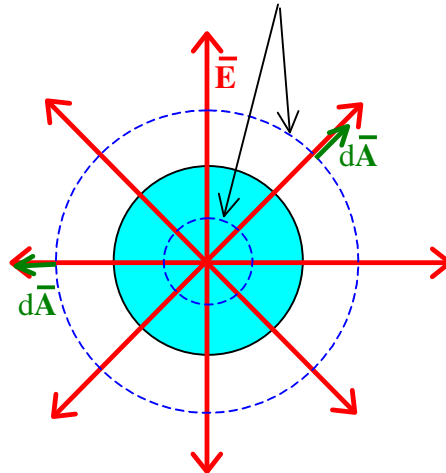
2. Best Gaussian surface: If we choose a sphere of radius r as the Gaussian surface, the field will be perpendicular to the surface at all points, so

$$\vec{E} \cdot d\vec{A} = EdA$$

Moreover, since, by symmetry, E can depend only on r and all points on the sphere have the same radius, E is constant over the spherical surface.

This applies to both of the Gaussian spheres.

Gaussian spheres of radius r ; one inside and one outside the sphere of charge



3. Find Φ : $\Phi = \oint \vec{E} \cdot d\vec{A} = \oint EdA = E \oint dA = E(\text{surface area of sphere}) = E(4\pi r^2)$

4. Find Q_{enc} :(i) $r < R$: Enclosed charge = (charge density)(volume of Gaussian Sphere)

$$Q_{\text{enc}} = \rho \frac{4\pi r^3}{3}$$

(ii) $r > R$: Enclosed charge = (charge density)(Volume of charged Sphere)

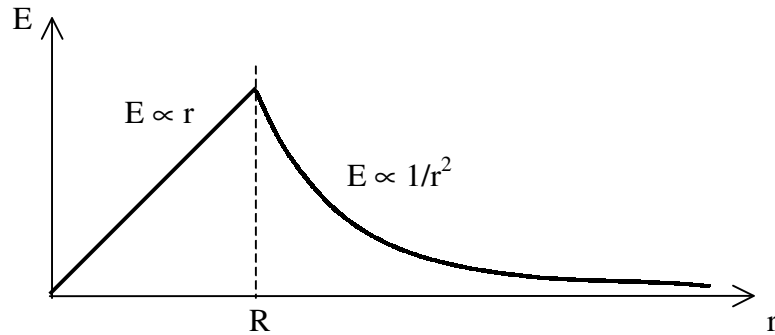
$$Q_{\text{enc}} = \rho \frac{4\pi R_1^3}{3}$$

5. Combine the results of (3) and (4) to find the electric field:

$$(i) \quad r < R: \quad 4\pi r^2 E = \frac{Q_{\text{enc}}}{\epsilon_0} = \rho \frac{4\pi r^3}{3\epsilon_0} \quad \Rightarrow \quad E = \frac{\rho r}{3\epsilon_0}.$$

$$(ii) \quad r > R: \quad 4\pi r^2 E = \frac{Q_{\text{enc}}}{\epsilon_0} = \rho \frac{4\pi R_1^3}{3\epsilon_0} \quad \Rightarrow \quad E = \frac{\rho R_1^3}{3\epsilon_0 r^2}.$$

Sketch of E vs. r:

(c) In the region between R_1 and R_2 , the field is given by

$$E = \frac{\rho r}{3\epsilon_0} - (E \text{ due to the sphere that has been removed})$$

$$\Rightarrow \quad E = \frac{\rho r}{3\epsilon_0} - \frac{\rho R_2^3}{3\epsilon_0 r^2} = \frac{\rho}{3\epsilon_0} \left[r - \frac{R_2^3}{r^2} \right].$$

For $r < R_2$, a Gaussian sphere would enclose no charge, so $E = 0$

$$\text{For } r > R_1, \text{ the field is } E = \frac{\rho(R_1^3 - R_2^3)}{3\epsilon_0 r^2}.$$

Answer 3

- (a) Relationship between electric field and potential: The potential difference between two points in space, a and b, is defined as

$$\Delta V = V_b - V_a = - \int_a^b \vec{E} \cdot d\vec{L}$$

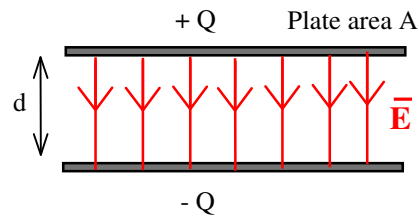
or: if $V(x,y,z)$ is the electric potential as a function of position in space, then

$$\vec{E}(x,y,z) = -\frac{\partial V}{\partial x} \hat{i} - \frac{\partial V}{\partial y} \hat{j} - \frac{\partial V}{\partial z} \hat{k}.$$

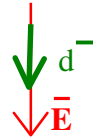
- (b) Assume that one plate carries charge $+Q$ and the other $-Q$. Ignoring edge effects, the electric field between the plates is uniform and given by $E = \sigma/K\epsilon_0$.

But $\sigma = Q/A$ so $E = \frac{Q}{K\epsilon_0 A}$

To find the potential difference between the plates, we integrate the electric field following any path from one plate to the other.



If we follow a field line, then $d\vec{L}$ is always parallel to \vec{E} and $\vec{E} \cdot d\vec{L} = EdL$



So $\Delta V = \int_0^d \frac{Q}{K\epsilon_0 A} dL = \frac{Qd}{K\epsilon_0 A}$. By definition, $C = \frac{Q}{\Delta V}$, so $C = \frac{K\epsilon_0 A}{d}$.

- (c) The total energy of the system is $U = \frac{1}{2} Q \Delta V = \frac{1}{2} \frac{Q^2 d}{K\epsilon_0 A}$ or

$$U = \frac{1}{2} C \Delta V^2 = \frac{1}{2} \frac{K\epsilon_0 A}{d} \frac{Q^2 d^2}{K^2 \epsilon_0^2 A^2} = \frac{1}{2} \frac{Q^2 d}{K\epsilon_0 A}.$$

Regard this energy as existing in the uniform electric field between the plates, which occupies a volume of Ad .

The energy per unit volume is therefore $u = \frac{1}{2} \frac{Q^2}{K\epsilon_0 A^2}.$

But $Q/A = K\epsilon_0 E$, so $u = \frac{1}{2} K\epsilon_0 E^2.$

(d) This system is equivalent to two capacitors in parallel.

For $x = a$, the capacitance is

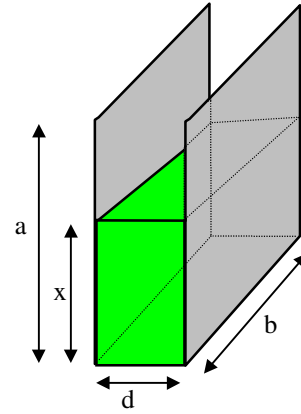
$$C_{\max} = \frac{4\epsilon_0 ab}{d}.$$

For $x < a$, the capacitance is

$$C(x) = \frac{4\epsilon_0 xb}{d} + \frac{\epsilon_0 (a-x)b}{d} = \frac{\epsilon_0 b}{d} (3x + a).$$

The ratio of the measured to the maximum value of capacitance is

$$\frac{C(x)}{C_{\max}} = R = \frac{3x + a}{4a} \quad \Rightarrow \quad x = \left[\frac{4R - 1}{3} \right] a.$$

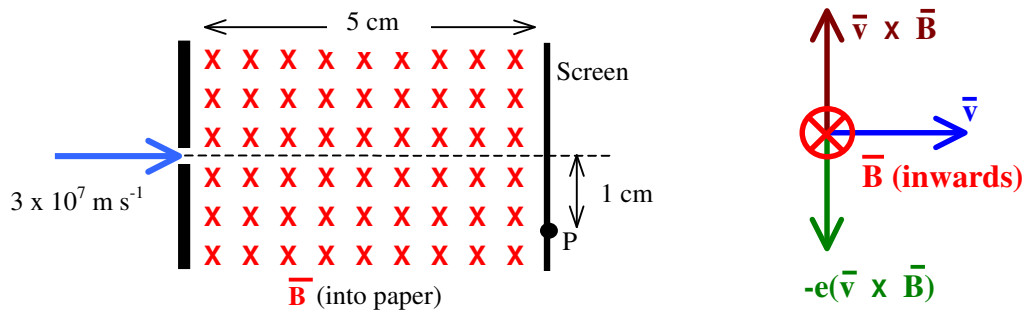


Answer 4

(a) $\vec{F} = Q[\vec{E} + (\vec{v} \times \vec{B})]$

The direction of the magnetic force is that of the cross product $\vec{v} \times \vec{B}$, which is perpendicular to both vectors. Since there is never any component of the magnetic force along the direction of \vec{v} , the speed of the particle cannot be changed by the magnetic force.

(b) The force on an electron moving in the magnetic field is $\vec{F} = -e(\vec{v} \times \vec{B})$



Now $\vec{v} \perp \vec{B}$ so $F = -e v B$

$\vec{F} \perp \vec{v}$ always and v never changes \Rightarrow circular motion

Putting the centripetal force equal to the magnetic force, we get

$$evB = \frac{m_e v^2}{r} \Rightarrow r = \frac{m_e v}{eB}.$$

$$e = 1.602 \times 10^{-19} \text{ C}$$

$$m_e = 9.11 \times 10^{-31} \text{ kg} \Rightarrow r = \frac{1.71 \times 10^{-4}}{B} \text{ m.}$$

$$v = 3.0 \times 10^7 \text{ m s}^{-1}$$

If the electrons strike the screen at point P, then

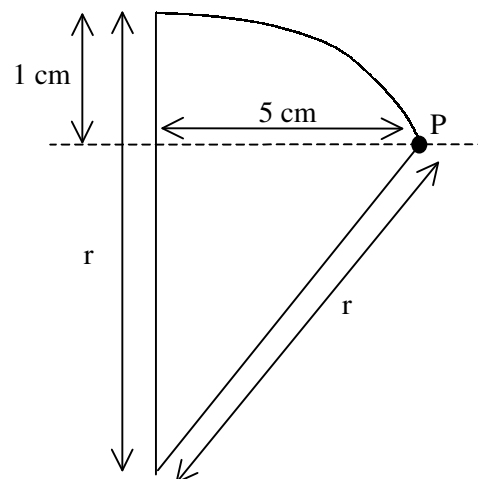
$$(r - 1)^2 + 5^2 = r^2.$$

$$\text{So } r^2 - 2r + 1 + 25 = r^2.$$

$$\Rightarrow r = 13 \text{ cm.}$$

$$\text{Therefore } 0.13 = \frac{1.71 \times 10^{-4}}{B}$$

$$\Rightarrow B = 1.31 \text{ mT.}$$



Answer 5

$$(a) \quad \oint \vec{B} \cdot d\vec{L} = \mu_0 I_{\text{enc}}.$$

\vec{B} = magnetic field

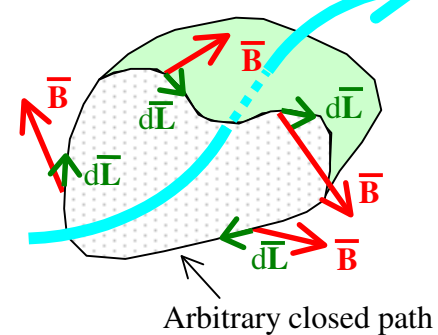
$d\vec{L}$ = an infinitesimal line element vector

μ_0 = permeability constant

I_{enc} = current flowing through a specified closed path

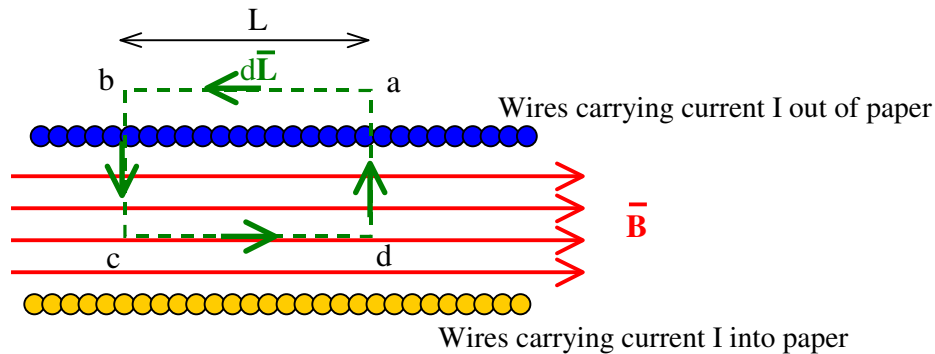
$\oint \Rightarrow$ The integral is to be carried out over a *closed* path

Total current I_{enc} passes through a bag-like surface whose opening is the path



Ampere' s law states that the line integral of the magnetic field around any closed path is equal to the total current flowing through the area defined by the path multiplied by μ_0 .

(b) Consider the solenoid seen as if it were sliced down the middle:



Apply Ampere' s law to the dotted rectangular path:

$$\oint \vec{B} \cdot d\vec{L} = \int_a^b \vec{B} \cdot d\vec{L} + \int_b^c \vec{B} \cdot d\vec{L} + \int_c^d \vec{B} \cdot d\vec{L} + \int_d^a \vec{B} \cdot d\vec{L}$$

$$\int_a^b \vec{B} \cdot d\vec{L} = 0 \quad \text{because } B = 0$$

$$\int_b^c \vec{B} \cdot d\vec{L} = \int_d^a \vec{B} \cdot d\vec{L} = 0 \quad \text{because } B = 0 \text{ outside the solenoid and } d\vec{L} \perp \vec{B} \text{ inside}$$

$$\int_c^d \vec{B} \cdot d\vec{L} = \int_c^d B dL = B \int_c^d dL \quad \text{as } d\vec{L} \text{ and } \vec{B} \text{ are parallel and } B \text{ is constant}$$

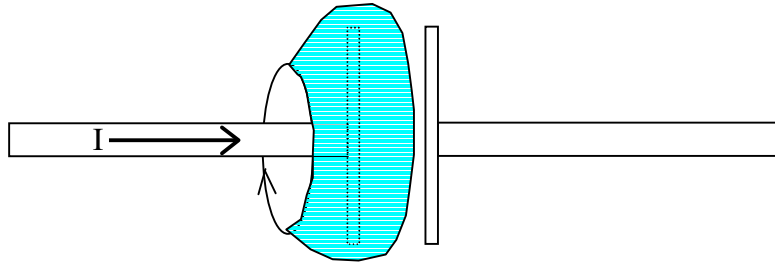
$$\int_c^d dL \text{ is just the length } L.$$

$$\text{Therefore } \oint \vec{B} \cdot d\vec{L} = BL$$

Enclosed current: nL wires are intercepted by the path, each carrying a current I . So $I_{\text{enc}} = nLI$.

Therefore $BL = \mu_0 nLI$ so $B = \mu_0 nI$.

(c) Consider a parallel plate capacitor being charged up by a current I :



We know that a moving charge produces a magnetic field, so the current in the wire will generate a magnetic field. Applying Ampere's law to the path shown, with the bag-like surface passing between the capacitor plates, we get

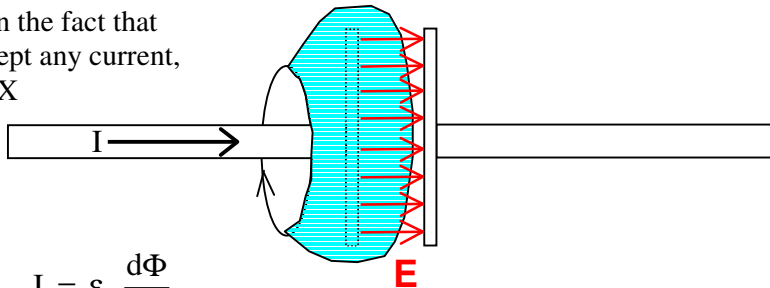
$$\oint \vec{B} \cdot d\vec{L} = \mu_0 I = 0 \Rightarrow \vec{B} \text{ must be zero, which we know cannot be true.}$$

Maxwell's modification is based on the fact that although the surface does not intercept any current, it DOES intercept ELECTRIC FLUX

How much flux is intercepted?

Let Q = charge on capacitor

$$\Phi = \frac{Q}{\epsilon_0} \Rightarrow \frac{d\Phi}{dt} = \frac{I}{\epsilon_0} \Rightarrow I = \epsilon_0 \frac{d\Phi}{dt}$$



\Rightarrow If we want $\oint \vec{B} \cdot d\vec{L} = \mu_0 I$ as for a straight wire, we must claim that

$$\oint \vec{B} \cdot d\vec{L} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi}{dt}$$

Term 1

Term 2

Continuous wire : Term 1 = $\mu_0 I$
Term 2 = 0

Capacitor : Term 1 = 0
Term 2 = $\mu_0 \epsilon_0 \frac{d\Phi}{dt} = \mu_0 I$

So the magnetic field is the same in each case

Maxwell showed that this is a general relation which holds always.

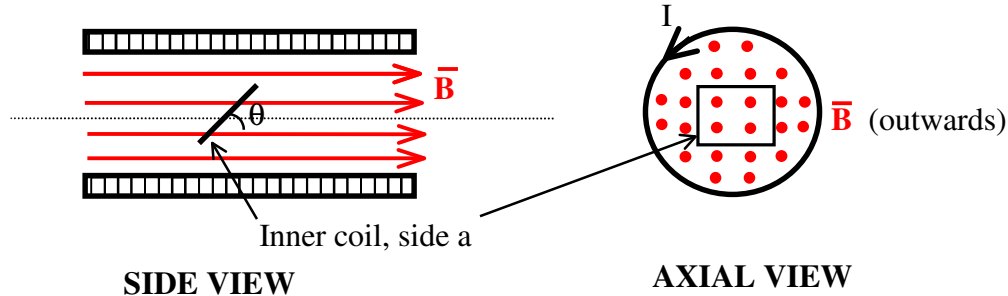
Answer 6

- (a) Mutual inductance between two circuits: if a circuit carries current I_1 and a second circuit links an amount of magnetic flux Ψ_{21} due to I_1 then the mutual inductance, M , is defined by

$$M = \Psi_{21}/I_1 \quad (\text{Inductance} = \text{Flux}/\text{Current}).$$

The emf induced in the second circuit is $E_2 = -M(dI_1/dt)$.

- (b) Long solenoid of radius a , n turns per unit length, current I_1 , $B = \mu_0 n I_1$.



$$(i) \quad M = \frac{\Psi_{21}}{I_1} = \frac{B(\text{No. of turns of inner coil})(\text{Projected area of inner coil})}{I_1} = \frac{\mu_0 n I_1 N a^2 \sin \theta}{I_1}.$$

Therefore $M = \mu_0 n N a^2 \sin \theta$.

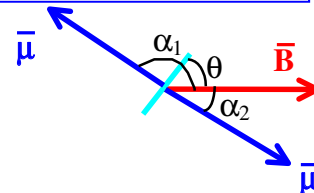
- (ii) Current in inner coil = I_2 .

\Rightarrow Magnetic dipole moment vector of inner coil has

Magnitude: $\mu = N a^2 I_2$ (current)(area)(no. of turns)

Direction: perpendicular to the coil, the exact direction depending on whether the current is clockwise or anticlockwise.

If I_2 is anticlockwise as seen looking in opposite direction to \vec{B}



If I_2 is anticlockwise as seen looking in opposite direction to \vec{B}

Magnitude of torque on inner coil = magnitude of $\vec{\tau} = \vec{\mu} \times \vec{B}$

$$\tau = \mu B \sin \alpha_1 \text{ or } \mu B \sin \alpha_2 \quad \Rightarrow \quad \tau = \mu B \cos \theta = (N a^2 I_2)(\mu_0 n I_1) \cos \theta.$$

$$\text{So } \tau = \mu_0 n N a^2 I_1 I_2 \cos \theta.$$

- (iii) Work done in rotating the coil through 180° starting from $\theta = 90^\circ$:

This corresponds to rotation between $\theta = \pi/2$ (the equilibrium position) and $\theta = 3\pi/2$. To overcome the magnetic torque, we must provide an external opposing torque $-\tau$. When the coil

is at an angle θ to the x axis and is rotated by an infinitesimal angle $d\theta$, the work done is thus $-\tau d\theta$. The total work done is therefore

$$W = - \int_{\pi/2}^{3\pi/2} \mu_0 n N a^2 I_1 I_2 \cos \theta d\theta = -\mu_0 n N a^2 I_1 I_2 [\sin \theta]_{\pi/2}^{3\pi/2} = -\mu_0 n N a^2 I_1 I_2 [-1 - 1] = 2\mu_0 n N a^2 I_1 I_2$$

The new position is one of unstable equilibrium: the torque increases as the coil moves out of position. If the coil were displaced slightly from this position, the torque would rotate it back towards the stable equilibrium position ($\theta = 90^\circ$).